## On the Navier–Stokes Equations with Coriolis Force Term

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I will talk on the Navier–Stokes equations in  $D \subset \mathbb{R}^3$ :

$$(NS) \begin{cases} u_t - \Delta u + (u \cdot \nabla)u + \Omega e_3 \times u + \nabla p = 0, \quad \nabla \cdot u = 0 \quad \text{in } D \times (0, T), \\ u_{\partial D} = 0, \quad u_{t=0} = u_0, \end{cases}$$

Here  $u = u(x, t) = (u_1, u_2, u_3)$  is an unknown velocity field and p = p(x, t) is an unknown pressure. In the equation (NS)  $e_3$  is the vertical unit vector (= (0, 0, 1)) and the term  $\Omega e_3 \times u$  is called the Coriolis force term with the Coriolis parameter  $\Omega \in \mathbb{R}$ .

The Coriolis force term is obtained in the following way. Consider the standard Navier–Stokes equation in  $\mathbb{R}^3$ :

$$v_t - \Delta v + (v \cdot \nabla)v + \nabla q = 0, \quad \nabla \cdot v = 0$$

with initial velocity

$$v|_{t=0}(y) = v_0(y) + (\Omega/2)e_3 \times y = v_0(y) + (\Omega/2)Jy \quad \text{with } J = \begin{pmatrix} 0 & -1 & 0 \\ 1 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Note that  $\operatorname{rot} v|_{t=0} = \operatorname{rot} v_0 + \Omega e_3$ . In this equation we introduce the transform

$$\begin{cases} u(x,t) = \exp\{-\Omega Jt/2\} v(y,t) - (\Omega/2) Jx \\ p(x,t) = \exp\{-\Omega Jt/2\} q(y,t) + \Omega^2 |y'|^2/8. \end{cases}$$

Here  $y = (y', y_3) = \exp \{\Omega J t/2\} x$ . Then  $(u, \nabla p)$  satisfies (NS).

From a meteorological point of view, if Coriolis parameter  $\Omega$  is sufficiently large, flow will be independent of vertical direction  $x_3$  asymptotically, that is 3 dim. flow will close to 2 dim. flow as  $|\Omega| \to \infty$ . This phenomena is called the Taylor– Proudman theorem. Our main purpose is to study this singular perturbation problem. In my talk I will discuss the following results.

For a fixed  $\Omega$ , we obtained several unique existence theorems of the Cauchy problem and boundary value problem  $(\mathbb{R}^3_+)$  in some function spaces which include periodic functions, almost periodic functions and some  $L^{\infty}$  functions.

In the case of the stationary problem in the Poincaré domain we noted some asymptotic behavior of velocity field in  $L^2(D)$  as  $\Omega \to \infty$ .

My talk is based on joint works with Prof. Y. Giga, Prof. K. Inui, Prof. A. Mahalov and Prof. J. Saal:

- Y. Giga, K. Inui, A. Mahalov and S. Matsui; Uniform local solvability for the Navier-Stokes equations with the Coriolis force, Methods and Applications of Analysis, Vol. 12 (2005) 381–384
- Y. Giga, K. Inui, A. Mahalov and S. Matsui; Navier-Stokes Equations in a Rotating Frame in ℝ<sup>3</sup> with Initial Data Nondecreasing at Infinity, Hokkaido Math. J., Vol. 35 (2006) 321–364.
- 3. Y. Giga, K. Inui, A. Mahalov, S. Matsui, and J. Saal; Rotating Navier-Stokes Equations in R<sup>3</sup><sub>+</sub> with Initial Data Nondecreasing at Infinity: The Ekman Boundary Layer Problem, to appear in Archive for Rational Mechanics and Analysis.