# Quenched Invariance Principle for a Reflecting Diffusion in a Continuum Percolation Cluster 

By Yutaka takeuchi

(Received Apr., 2022)
(Revised , )


#### Abstract

We consider a continuum percolation built over stationary ergodic point processes. Assuming that the occupied region has a unique unbounded cluster and the cluster satisfies volume regularity and isoperimetric condition, we prove a quenched invariance principle for reflecting diffusions on the cluster.


## 1. Introduction

### 1.1. Models and result

Let $\Omega$ be a configuration space on $\mathbb{R}^{d}$, namely

$$
\Omega=\left\{\sum_{i=1}^{\infty} \delta_{x_{i}} \mid \sum_{i=1}^{\infty} \delta_{x_{i}}(K)<\infty \text { for all compact set } K\right\}
$$

Here $\delta_{x}$ denotes the Dirac measure. $\Omega$ is equipped with the $\sigma$-field $\mathcal{B}(\Omega)$ which is generated by the sets $\{\omega \in \Omega \mid \omega(A)=n\}, A \in \mathcal{B}\left(\mathbb{R}^{d}\right), n \in \mathbb{N}$. We identify each configuration $\omega=\sum_{i=1}^{\infty} \delta_{x_{i}} \in \Omega$ with the subset $\left\{x_{i}\right\}_{i}$ of $\mathbb{R}^{d}$ if $\left\{x_{i}\right\}_{i}$ are distinct. Fix $\rho>0$. For $\omega \in \Omega$, define the subset $L(\omega)$ by

$$
L(\omega)=\bigcup_{x \in \omega} B(x, \rho),
$$

where $B(x, \rho)$ is the Euclidean open ball with center $x$ and radius $\rho$. Let $W(\omega)$ be the unbounded connected component of $L(\omega)$ if there is a unique unbounded component. Otherwise, we set $W(\omega)=\varnothing$ by convention. We call $W(\omega)$ the continuum percolation cluster. This is the continuum analogue of the discrete site percolation cluster. By definition, $W(\omega)$ can be written as $\bigcup_{x \in I(\omega)} B(x, \rho)$ with some subset $I(\omega)$ of $\omega$. We also take $\rho^{\prime} \geq \rho$ and we introduce the modified cluster $W^{\prime}(\omega)$ given by

$$
W^{\prime}(\omega) \equiv W_{\rho^{\prime}}(\omega)=\bigcup_{x \in I(\omega)} B\left(x, \rho^{\prime}\right) .
$$

Set

$$
\Delta=\left\{\begin{array}{l|l}
\omega \in \Omega & \begin{array}{l}
\omega=\sum_{i} \delta_{x_{i}}, \quad|x-y|=2 \rho(\text { for some } x, y \in \omega) \\
\text { and each } x_{i} \text { is distinct }
\end{array}
\end{array}\right\}
$$

[^0]Define the subset $\widehat{\Omega}$ of $\Omega$ by

$$
\widehat{\Omega}=\left\{\omega \in \Omega \backslash \Delta \mid 0 \in W^{\prime}(\omega)\right\} .
$$

We define the shift $\tau_{z}: \Omega \rightarrow \Omega$ by

$$
\tau_{z} \omega(A)=\omega(A+z)=\sum_{x \in \omega} \delta_{x}(A+z)=\sum_{x \in \omega} \delta_{x-z}(A)
$$

where $A+z=\{x+z \mid x \in A\}$ for $z \in \mathbb{R}^{d}$.
Let $\mathbb{P}$ be a probability measure on $\Omega$. It is called a point process on $\mathbb{R}^{d}$.
Assumption 1. Assume that $\mathbb{P}$ satisfies the following conditions:
(1) $\mathbb{P}$ is stationary and ergodic with respect to $\left\{\tau_{x}\right\}_{x}$.
(2) $\mathbb{P}(\widehat{\Omega})>0$ and $\mathbb{P}(\Delta)=0$.

Set $\widehat{\mathbb{P}}(\cdot)=\mathbb{P}(\cdot \cap \widehat{\Omega}) / \mathbb{P}(\widehat{\Omega})$ and we denote its expectation by $\widehat{\mathbb{E}}$. Note that $W^{\prime}(\omega)$ is a Lipschitz domain for $\widehat{\mathbb{P}}$-a.e. $\omega \in \widehat{\Omega}$. Thanks to $\mathbb{P}(\Delta)=0$, we only need to consider whether $0 \in W(\omega)$ or not. The existence of a unique unbounded component is an important problem in the study of continuum percolation theory and there are many studies. If $\mathbb{P}$ is a Poisson point process and the radius $\rho$ is bigger than the critical value, it is known that there is a unique unbounded component (see [25]).

We denote the Euclidean inner product by $\langle\cdot, \cdot\rangle$. Let $a: \Omega \rightarrow \mathbb{R}^{d \times d}$ be a positivedefinite symmetric matrix. Define the bilinear form $\mathcal{E}$ on $L^{2}\left(W^{\prime}(\omega), d x\right)$ by

$$
\begin{equation*}
\mathcal{E}^{\omega}(u, v)=\int_{W^{\prime}(\omega)}\left\langle a\left(\tau_{x} \omega\right) \nabla u(x), \nabla v(x)\right\rangle d x \tag{1.1}
\end{equation*}
$$

Let $\mathcal{F}^{\omega}$ be the completion of $C_{c}^{\infty}(\bar{W})$ with respect to $\mathcal{E}(\cdot, \cdot)+\|\cdot\|_{L^{2}\left(W^{\prime}(\omega)\right)}$.
Assumption 2. There exist constants $\lambda, \Lambda>0$ such that

$$
\lambda|\xi|^{2} \leq\langle a(\omega) \xi, \xi\rangle \leq \Lambda|\xi|^{2}
$$

holds for all $\xi \in \mathbb{R}^{d}$ and $\widehat{\mathbb{P}}$-almost all $\omega$.
According to [17] and [18], under Assumption 2, we have that the Dirichlet form $\left(\mathcal{E}^{\omega}, \mathcal{F}^{\omega}\right)$ is strongly local and regular $\widehat{\mathbb{P}}$-almost surely. Hence, we have the associated conservative diffusion $X_{t}^{\omega}$. It is called a reflecting diffusion since the domain $\mathcal{F}^{\omega}$ of Dirichlet form corresponds to a reflecting boundary condition (more precisely, the Neumann boundary condition).

We further impose an assumption for reflecting diffusions:
Assumption 3. The reflecting diffusion $X_{t}^{\omega}$ has a transition density $p_{t}^{\omega}(\cdot, \cdot)$ for $\widehat{\mathbb{P}}$-almost all $\omega$.

When we consider the case that $a(\omega)=1 / 2 I_{d}, X_{t}$ is the reflecting Brownian motion. Bass and Hsu ([6]) study a condition under which $X_{t}$ has a density. Recently, Matsuura ([24]) considered more detailed conditions. Our model satisfies the conditions of

Matsuura [24], and hence the reflecting Brownian motion has a density.
Next, we impose a geometric condition that plays an important role. Let $x \in \mathbb{R}^{d}$ and $R>0$. Let $W^{\prime}(\omega, x, R)$ be a connected component of $W^{\prime}(\omega) \cap B(x, R)$ containing $x$. Throughout the paper, we write $W^{\prime}=W^{\prime}(\omega)$ and $W_{R}^{\prime}=W^{\prime}(\omega, 0, R)$ if there is no confusion.

Assumption 4 (VOLUME REGULARITY AND ISOPERIMETRIC CONDITION).
(1) For $\widehat{\mathbb{P}}$-almost every $\omega$ there exists positive constant $C_{V}$ such that for a.e. $x \in W^{\prime}(\omega)$

$$
C_{V} R^{d} \leq\left|W^{\prime}(\omega, x, R)\right|, \quad R \geq R_{V}
$$

holds for some positive constant $R_{V}$, where $|\cdot|$ denotes the Lebesgue measure.
(2) There exists $\theta \in(0,1)$ such that for $\widehat{\mathbb{P}}$-almost all $\omega \in \widehat{\Omega}$, there exist $R_{I}>0$ and $c_{H}>0$ such that

$$
C_{\mathrm{IL}}:=\inf \left\{\frac{\mathcal{H}_{d-1}\left(W^{\prime} \cap \partial O\right)}{\left|W^{\prime} \cap O\right|^{\frac{d-1}{d}}} \left\lvert\, \begin{array}{l}
O \subset B(0, R) \text { is connected open, }  \tag{1.2}\\
R \geq R_{I},|O| \geq R^{\theta}
\end{array}\right.\right\}>0
$$

and

$$
C_{\mathrm{IS}}:=\inf \left\{\frac{\mathcal{H}_{d-1}\left(W^{\prime} \cap \partial O\right)}{\left|W^{\prime} \cap O\right|^{\frac{d-1}{d}}} \left\lvert\, \begin{array}{l}
O \text { is bounded open, }  \tag{1.3}\\
\mathcal{H}_{d-1}\left(W^{\prime} \cap \partial O\right)<c_{H}
\end{array}\right.\right\}>0
$$

hold, where $\mathcal{H}_{d-1}$ denotes the $(d-1)$-dimensional Hausdorff measure.
The main result of this paper is the following.
Theorem 1.1 (Quenched invariance principle). Let $d \geq 2$ and $\rho^{\prime} \geq \rho>0$. Assume that Assumptions 1, 2, 3 and 4 hold. Let $P_{0}^{\omega}$ be the law of $\left\{X_{t}^{\omega}\right\}_{t}$ starting at 0 . Then for $\widehat{\mathbb{P}}$-a.s. $\omega$, the scaled process $\left\{\varepsilon X_{\varepsilon^{-2} t}^{\omega}\right\}_{t}$ under $P_{0}^{\omega}$ converges in law to a Brownian motion with non-degenerate covariance matrix as $\varepsilon$ tends to 0 .

REmARK 1.2. When $\rho^{\prime}>\rho$, the condition (1.3) is automatically satisfied. On the other hand, the condition (1.3) doesn't hold when $\rho^{\prime}=\rho$ in general.

Example 1.3 (Reflecting Brownian motion on Poisson Boolean model). Let $\mathbb{P}$ be a Poisson point process with intensity $\lambda$. We take the radius $\rho$ greater than the critical radius $\rho_{c}$. We also take the radius $\rho^{\prime}>\rho$. Set $a(\omega)=1 / 2 I_{d}$. Then the corresponding diffusion is the reflecting Brownian motion. We can easily verify that Assumption 1 and 2 are satisfied. Assumption 3 is satisfied by the result in [24]. Now we need to check the Assumption 4. We check this by comparing the modified cluster $W^{\prime}(\omega)$ to the Bernoulli site percolation model as in [33]. To do this, we prepare a percolation model on $\delta \mathbb{Z}^{d}, \delta>0$. We call the model the $\delta$-approximating Boolean model. Set $G(z, \delta)=[-\delta / 2, \delta / 2]^{d}+z$. Take $\delta \in(0,1)$. We say that a site $z \in \delta \mathbb{Z}^{d}$ is $\delta$-open if $G(z, \delta) \subset W^{\prime}(\omega)$. Let $V_{\delta}$ be the collection of $\delta$-open sites on $\delta \mathbb{Z}^{d}$. We introduce the graph $\mathcal{G}_{\delta}=\left(V_{\delta}, E_{\delta}\right)$ with vertex set $V_{\delta}$ and edge set $E_{\delta}=\left\{\left\{z, z^{\prime}\right\} \subset V_{\delta}| | z-z^{\prime} \mid=\delta\right\}$. For $x \in V_{\delta}$, let $B_{\delta}^{\omega}(x, R)$ be an open ball of $V_{\delta}$ with respect to the graph distance. For
a cube $Q=G(z, n \delta)$, write $Q^{+}=G\left(z, \frac{3}{2} n \delta\right)$. We say that a cluster $\mathcal{C}$ in a cube $Q$ is crossing for a cube $Q^{\prime} \subset Q$ if for all $d$-directions there exists an open path in $\mathcal{C} \cap Q^{\prime}$ that connects the two opposing faces of $Q^{\prime}$. Let $\mathcal{C}^{\vee}(Q)$ be the largest cluster in $Q$. As in [4], set

$$
\begin{aligned}
R_{0}(Q)=\{ & \text { there exists a unique crossing cluster } \mathcal{C} \text { in } Q^{+} \text {for } Q^{+} \text {such that } \\
& \text { all open paths contained in } Q^{+} \text {of diameter greater than } n \delta / 8 \\
& \text { are connected to } \mathcal{C} \text { in } Q^{+} \text {and } \mathcal{C} \text { is crossing for each cube } \\
& \text { contained in } Q \text { whose side length is greater than or equal to } n \delta / 8\}
\end{aligned}
$$

and

$$
R(Q)=R_{0}(Q) \cap\left\{\mathcal{C}^{\vee}(Q) \text { is crossing for } Q\right\} \cap\left\{\mathcal{C}^{\vee}\left(Q^{+}\right) \text {is crossing for } Q^{+}\right\}
$$

First we consider the following estimate:

$$
\begin{equation*}
\mathbb{P}\left(R(G(z, k \delta))^{c}\right) \leq c \exp \left(-c^{\prime} k\right), k \in \mathbb{N} \tag{1.4}
\end{equation*}
$$

In [4, Lemma 2.8], the estimate (1.4) is proved in the case of the Bernoulli site percolation. We can generalize the estimate (1.4) to the $\delta$-approximating Boolean model by renormalization used in [33]. Hence as in [4], we can show that the process $\left\{1_{R(G(z, k \delta))}\right\}_{z \in \delta \mathbb{Z}^{d}}$ dominates the Bernoulli site percolation with parameter $q_{\delta}^{*}(k)$, which tends to 1 as $k \rightarrow \infty$ and satisfies $q_{\delta}^{*}(k) \leq q_{\delta^{\prime}}^{*}(k)$ for $\delta \geq \delta^{\prime}$. Therefore, for sufficiently large $n$ and a cube $Q=G\left(z_{0}, n \delta\right)$, we can prove the weak relative isoperimetric inequality

$$
\begin{equation*}
\#\left\{\left\{z, z^{\prime}\right\}\left|z \in A, z^{\prime} \in \mathcal{C}^{\vee}(Q)-A,\left|z-z^{\prime}\right|=\delta\right\} \geq c_{1} n^{-1} \# A\right. \tag{1.5}
\end{equation*}
$$

holds for connected subset $A \subset \mathcal{C}^{\vee}(Q)$ with $\# A \leq 1 / 2 \# \mathcal{C}^{\vee}(Q)$ such that $\mathcal{C}^{\vee}(Q)-A$ is also connected, and the volume regularity

$$
\begin{equation*}
c_{2} R^{d} \leq \# B_{\delta}^{\omega}(z, R) \leq c_{2}^{\prime} R^{d} \tag{1.6}
\end{equation*}
$$

holds for all $z \in \mathcal{C}^{\vee}\left(Q^{+}\right) \cap G\left(z_{0}, 5 / 6 n \delta\right)$ with $G(z, R+k)^{+} \subset Q^{+}$and $R \in\left(c_{H} n^{\alpha}, n\right)$ as in [4, Proposition 2.11 and Theorem 2.23]. Then by [10, Lemma 2.10], there exist $\theta \in(0,1)$ and $N_{3}$ such that $V_{\delta}$ satisfies the isoperimetric inequality for large sets

$$
\begin{equation*}
\frac{\#\left\{\left\{z, z^{\prime}\right\} \in E_{\delta} \mid z \in A, z^{\prime} \in V_{\delta}-A\right\}}{(\# A)^{\frac{d-1}{d}}} \geq c_{3} \tag{1.7}
\end{equation*}
$$

holds for $R \geq N_{3}$ and $A \subset B_{\delta}^{\omega}(0, R)$ with $\# A \geq R^{\theta}$. Note that $B_{\delta}^{\omega}(0, R)$ is contained in the unique infinite cluster of $V_{\delta}$. Furthermore, we can take constants $c_{1}, c_{2}, c_{2}^{\prime}, c_{3}>0$, and $N_{3}>0$ independently of $\delta \in\left(0, \delta_{0}\right)$ for some $\delta_{0}$ (This is because these constants depend only on $k$ and $q_{\delta}^{*}(k)$ is decreasing in $\delta$.).

Now we can check Assumption 4. Since $\rho^{\prime}>\rho$, the condition (1.3) is satisfied (see Remark 1.2.). Because $B_{\delta}^{\omega}(x, R / \delta)$ is contained in $W^{\prime}(\omega, x, R)$, we have $\left|W^{\prime}(\omega, x, R)\right| \geq$ $\delta^{d} \#\left|B_{\delta}^{\omega}(x, R / \delta)\right|$. Hence we can easily see that (1) of Assumption 4 holds. To check (2)
of Assumption 4, we introduce some notations. Set

$$
\mathbb{G}_{\delta}=\left\{G(z, \delta) \mid z \in \mathbb{Q}^{d}\right\}, \quad \mathbb{G}=\bigcup_{\delta \in(0,1) \cap \mathbb{Q}} \mathbb{G}_{\delta}, \quad \mathcal{Q}_{\mathrm{fin}}=\left\{\bigcup_{i=1}^{N} Q_{i} \mid Q_{i} \in \mathbb{G}\right\}
$$

Take a connected subset $D \in \mathcal{Q}_{\text {fin }}$ with $W^{\prime}(\omega) \cap D \neq \varnothing$. We can choose $\delta>0$ so that $D$ is written as $\bigcup_{i=1}^{N} G\left(z_{i}, \delta\right), z_{i} \in \delta \mathbb{Z}^{d}, i=1, \ldots, N$. Moreover, we can take $\delta>0$ smaller if we need. Set $Z=\left\{z_{1}, \ldots, z_{N}\right\}$. Let $Z_{o}$ be the collection of $\delta$-open sites of $Z$ and $\partial_{\text {ext }} Z_{o}$ be the external boundary of $Z_{o}$. Observe that at least one face of cubes in $\bigcup_{z \in \partial_{\text {ext }} Z_{o}} G(z, \delta)$ belongs to $W^{\prime}(\omega) \cap \partial D$. Therefore we have $\mathcal{H}_{d-1}\left(W^{\prime}(\omega) \cap \partial Q\right) \geq$ $\delta^{d-1} \# \partial_{\text {ext }} Z_{o}$. Next, let $Z_{c}=\left\{z \in Z \backslash Z_{o} \mid G(z, \delta) \cap W^{\prime}(\omega) \neq \varnothing\right\}$. Then we have that $\left|W^{\prime}(\omega) \cap D\right| \leq \delta^{d}\left(\# Z_{o}+\# Z_{c}\right) \leq 2 \delta^{d} \# Z_{o}$ for sufficiently small $\delta$. Hence if $|D| \geq R^{\theta}$, using the isoperimetric inequality (1.7), we have that $\mathcal{H}_{d-1}\left(W^{\prime}(\omega) \cap \partial D\right) \geq C_{1}\left|W^{\prime}(\omega) \cap D\right|^{\frac{d-1}{d}}$. For a general connected open subset $O \subset B(0, R)$, we can take a sequence of $D_{i} \in \mathcal{Q}_{\text {fin }}$ such that $\left|W^{\prime}(\omega) \cap D_{i}\right| \rightarrow\left|W^{\prime}(\omega) \cap O\right|$ and $\mathcal{H}_{d-1}\left(W^{\prime}(\omega) \cap \partial D_{i}\right) \rightarrow \mathcal{H}_{d-1}\left(W^{\prime}(\omega) \cap \partial O\right)$. Hence, the condition (1.2) holds.

Example 1.4. Let $\left\{Z_{z}\right\}_{z \in \mathbb{Z}^{d}}$ be a collection of i.i.d. Bernoulli random variables with parameter $p \in(0,1)$ and $U$ be the uniform random variable on $[0,1)^{d}$. Let $\mathbb{P}$ be the distribution of $\sum_{z \in \mathbb{Z}} Z_{z} \delta_{z+U}$. Let $\rho_{c}>0$ be the critical radius. We remark that $\rho_{c}=1 / 2$ if the parameter $p$ is greater than the critical value of the Bernoulli site percolation. Take $\rho>\rho_{c}$ and let $\rho^{\prime}=\rho$. Set $a(\omega)=1 / 2 I_{d}$. Then we can check Assumption 1, 2, and 3 hold. Similarly to Example 1.3, we can check (1) of Assumption 4 and (1.2) hold. To check (1.3), observe that by construction, there is only a finite number of possible distances between overlapping balls. Indeed, their values are of the form $|x|, x \in \mathbb{Z}^{d}$, and the maximum is strictly less than $2 \rho$. Therefore, computing $\frac{\mathcal{H}_{d-1}\left(W^{\prime} \cap \partial O\right)}{\left|W^{\prime} \cap O\right|^{\frac{d-1}{d}}}$ in each case, we can check (1.3).

Early in the history of homogenization, Kipnis and Varadhan[21] proved that the annealed invariance principle for random walk on supercritical (bond) percolation cluster. After two decades, Sidoravicius and Sznitman [31] proved the quenched invariance principle for random walk on supercritical percolation cluster when the dimension is more than or equal to four. In 2007, Berger and Biskup [7] proved quenched invariance principle for random walk on the supercritical percolation cluster including two and three dimensions. (Note that one-dimensional percolation cluster is infinite only when the probability that bond is open is equal to one and in that case the QIP follows from classical Donsker's invariance principle.) Mathieu and Piatnitski [23] gave another proof. Quenched invariance principle for the random conductance model, which is a more general model, was shown by many authors ([1], [2], [5], [8], [22]). Recently, the quenched invariance principle for random conductance model with more general assumptions was shown ([3], [10]). They include not only the bond percolation cluster but also percolation clusters in models with long-range correlations and random conductance models defined by level sets of the Gaussian free field. In these papers, the proofs are based on an analysis on (general) weighted graph. They also use geometric conditions such as a relative isoperimetric inequality and a volume regularity.

In these results, they consider discrete settings. When we consider the continuum
settings, the diffusion process is one of the fundamental objects. For the diffusion process, there are many homogenization results. In [29] and [30], annealed invariance principles were shown. (precisely, slightly stronger results were proved but not quenched results.) For quenched results, there are [9], [13], [14],[12], [26] and [32].

The homogenization problem of the continuum percolation cluster is considered by some researchers. Tanemura [34] and Osada ([27], [28]) proved the annealed invariance principle for the reflecting Brownian motion in the continuum percolation cluster. However, to our knowledge, there are no quenched results.

### 1.2. Method

A basic but powerful method to prove a quenched invariance principle is harmonic embedding. A key ingredient of this approach is the corrector, a random function, $\chi: \mathbb{R}^{d} \times$ $\Omega \rightarrow \mathbb{R}^{d}$ which is the solution of Poisson type equation

$$
\begin{equation*}
\mathcal{L}^{\omega} \chi^{k}(x, \omega)=\mathcal{L}^{\omega} \Pi^{k}(x),(k=1, \ldots, d) \tag{1.8}
\end{equation*}
$$

where $\mathcal{L}^{\omega}$ is the corresponding generator of the diffusion $\left\{X_{t}^{\omega}\right\}_{t}$ and $\Pi^{k}(x)=x_{k}$ is the projection to the $k$-th coordinate. Then $y^{k}=\Pi^{k}-\chi^{k}$ is a harmonic function. This implies that

$$
M_{t}^{\omega}=X_{t}^{\omega}-\chi\left(X_{t}^{\omega}, \omega\right)
$$

is a martingale and a quenched invariance principle for the martingale part $M$ can be easily shown by standard arguments. In order to obtain a quenched invariance principle for the process $X$, it suffices to show that for any $T>0$ and $\widehat{\mathbb{P}}$-a.s. $\omega$

$$
\lim _{\varepsilon \rightarrow 0} \sup _{0 \leq t \leq T} \varepsilon\left|\chi\left(X_{t / \varepsilon^{2}}^{\omega}, \omega\right)\right|=0 \quad \text { in } P_{0}^{\omega} \text { - probability, }
$$

which can be deduced from the $L^{\infty}$-sublinearity of the corrector:

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \sup _{x \in W_{R}^{\prime}}\left|\chi_{\varepsilon}(x, \omega)\right|=0, \quad \widehat{\mathbb{P}} \text {-a.s. } \tag{1.9}
\end{equation*}
$$

where $\chi_{\varepsilon}(x, \omega)=\varepsilon \chi(x / \varepsilon, \omega)$. The non-degeneracy of the covariance matrix follows from the $L^{\infty}$-sublinearity and the ergodic theorem. So, important things are the followings:
(1) How to construct the corrector.
(2) How to prove the $L^{\infty}$-sublinearity.

One way to construct the corrector is to decompose the space of random function into a space of "potential" and its "orthogonal" space.

Although the $L^{\infty}$-sublinearity is difficult to show, a $L^{p}$-sublinearity, which is the equation obtained from (1.9) replacing the $L^{\infty}$-norm by the $L^{p}$-norm, is easily shown by the ergodic theorem. In [9] and [10], it is mentioned that the $L^{\infty}$-sublinearity follows from the $L^{p}$-sublinearity and a maximal inequality (3.21). They also suggested in these papers that Moser's iteration scheme is useful to obtain the maximal inequality. Since this method is analytic and robust, we can use this for the reflecting diffusion case.

Throughout the proof, we will mainly follow the argument that appeared in Chiarini and Deuschel [9]. However, different from their paper, we need to consider the boundary
effect. For example, we have to solve the equation (1.8) with the Neumann boundary condition. Due to this boundary condition, the proof of the maximal inequality (3.21) is more complicated. To consider this boundary condition, we have to analyze the space $\mathcal{F}^{\omega}$. Unlike the space $H_{0}^{1}(W)$ (the closure of $C_{c}^{\infty}(W)$ with respect to a Sobolev norm), the trace of functions that belong to $\mathcal{F}^{\omega}$ doesn't vanish in general. Thus, we have to control the boundary effect to obtain the maximal inequality. To overcome this difficulty, we kill the process $X_{t}^{\omega}$ at the boundary $\partial W^{\prime}(\omega, x, R)-\partial W^{\prime}(\omega)$ and consider local inequalities in Section 3.

## 2. Harmonic embedding

For $x \in \mathbb{R}^{d}$ and $F: \Omega \rightarrow \mathbb{R}^{d}$, define $T_{x} F$ by $\left(T_{x} F\right)(\omega)=F\left(\tau_{x} \omega\right)$. Set
$\widehat{\mathbb{L}}^{2}=\left\{F: \widehat{\Omega} \rightarrow \mathbb{R}^{d} \mid \widehat{\mathbb{E}}[\langle a F, F\rangle]<\infty\right\}$. We endow it with inner product $\langle a F, G\rangle$ for $F, G \in \widehat{\mathbb{L}}^{2}$.

We define the $i$-th derivative $D_{i}$ by

$$
D_{i} U=L^{2}-\lim _{h \rightarrow 0} \frac{T_{h e_{i}} U-U}{h},
$$

where $e_{i}$ is the $i$-th coordinate vector and $U \in L^{2}(\widehat{\Omega}, \widehat{\mathbb{P}})$.
Denote the domain of $D_{i}$ by $\mathcal{D}\left(D_{i}\right)=\left\{U \in L^{2}(\widehat{\Omega}, \widehat{\mathbb{P}}) \mid D_{i} U\right.$ exists $\}$. Set

$$
\mathcal{C}=\left\{\int_{\mathbb{R}^{d}} f\left(\tau_{x} \omega\right) \varphi(x) d x \mid f \in L^{\infty}(\Omega, \widehat{\mathbb{P}}), \varphi \in C_{c}^{\infty}\left(\mathbb{R}^{d}\right)\right\}
$$

We can easily show that if $v(\omega)=\int_{\mathbb{R}^{d}} f\left(\tau_{x} \omega\right) \varphi(x) d x \in \mathcal{C}$, then the $i$-th derivative of $v$ exists and it is given by

$$
\begin{equation*}
D_{i} v(\omega)=-\int_{\mathbb{R}^{d}} f\left(\tau_{x} \omega\right) \partial_{i} \varphi(x) d x \tag{2.1}
\end{equation*}
$$

Moreover, if $v \in \mathcal{C}$, then $v \in \bigcap_{i=1}^{d} \mathcal{D}\left(D_{i}\right)$.
Definition 2.1. We define the gradient and the space of potentials as follows:
(1) For $v \in \mathcal{C}$, set $D v=\left(D_{1} v, \ldots, D_{d} v\right) \in \widehat{\mathbb{L}}^{2}$.
(2) Define the subspace $\widehat{\mathbb{L}}_{\text {pot }}^{2}$ of $\widehat{\mathbb{L}}^{2}$ by the closure in $\widehat{\mathbb{L}}^{2}$ of the set $\{D v \mid v \in \mathcal{C}\}$.

REmARK 2.2. (1) In the settings of [9], the mean of a potential is zero and the authors of [9] used this property to prove the sublinearity of the corrector. Remind that the space $\mathcal{C}$ and the space of potentials can be generalized on $(\Omega, \mathbb{P})$. Then, for all generalized potential $\tilde{U}$, we have that $\mathbb{E}[\tilde{U}]=0$. However, the equality $\widehat{\mathbb{E}}[U]=0$ does not hold for a general potential $U \in \widehat{\mathbb{L}}_{\text {pot }}^{2}$.
(2) For $v \in \mathcal{C}$, set $v(x, \omega)=v\left(\tau_{x} \omega\right)$. Then we have $D v(x, \omega)=\nabla v(x, \omega)$.

Lemma 2.3. Let $U \in \widehat{\mathbb{L}}_{\text {pot }}^{2}$. Then for all $\eta \in C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$ and $i, j=1, \ldots, d$, we have

$$
\begin{equation*}
\int_{\mathbb{R}^{d}} U_{i}\left(\tau_{x} \omega\right) \partial_{j} \eta(x) d x=\int_{\mathbb{R}^{d}} U_{j}\left(\tau_{x} \omega\right) \partial_{i} \eta(x) d x \tag{2.2}
\end{equation*}
$$

for $\widehat{\mathbb{P}}$-a.e. $\omega \in \Omega$.
Proof. We first consider the case $v \in \mathcal{C}$. Then $x \mapsto v\left(\tau_{x} \omega\right)$ is infinitely many times differentiable, $\widehat{\mathbb{P}}$-a.s. Integrating by parts we get

$$
\begin{aligned}
\int_{\mathbb{R}^{d}} D_{i} v\left(\tau_{x} \omega\right) \partial_{j} \eta(x) d x & =-\int_{\mathbb{R}^{d}} v\left(\tau_{x} \omega\right) \partial_{i} \partial_{j} \eta(x) d x \\
& =-\int_{\mathbb{R}^{d}} v\left(\tau_{x} \omega\right) \partial_{j} \partial_{i} \eta(x) d x \\
& =\int_{\mathbb{R}^{d}} D_{j} v\left(\tau_{x} \omega\right) \partial_{i} \eta(x) d x
\end{aligned}
$$

For general $U \in \widehat{\mathbb{L}}_{\text {pot }}^{2}$ take approximations and use the fact that as $n \rightarrow \infty, \nabla v_{n} \rightarrow U$ in $\widehat{\mathbb{L}}^{2}$ implies $D_{i} v_{n}(\tau . \omega) \rightarrow U_{i}(\tau . \omega)$ in $L_{\text {loc }}^{1} \widehat{\mathbb{P}}$-a.s.

Let $\pi^{k}$ be the unit vector in the $k$-th direction. Since $\pi^{k} \in \widehat{\mathbb{L}}^{2}$, for each $k=1, \ldots, d$, there exist functions $U^{k} \in \widehat{\mathbb{L}}_{\text {pot }}^{2}$ and $R^{k} \in\left(\widehat{\mathbb{L}}_{\text {pot }}^{2}\right)^{\perp}$ such that $\pi^{k}=U^{k}+R^{k}$. Then we build the corrector starting from the functions $U^{k} \in \widehat{\mathbb{L}}_{\text {pot }}^{2}$. For $k=1, \ldots, d$, we define the corrector to be the function $\chi^{k}: \mathbb{R}^{d} \times \Omega \rightarrow \mathbb{R}$ such that

$$
\chi^{k}(x, \omega)=\sum_{j=1}^{d} \int_{0}^{1} x_{j} U_{j}^{k}\left(\tau_{t x} \omega\right) d t
$$

Proposition 2.4 (Weak differentiability). For $k=1, \ldots d$, the function $x \mapsto$ $\chi^{k}(x, \omega)$ is in $L_{\text {loc }}^{1}\left(\mathbb{R}^{d}\right)$, weakly differentiable $\mathbb{P}$-a.s. and $\partial_{i} \chi^{k}(x, \omega)=U_{i}^{k}\left(\tau_{x} \omega\right)$ for $i=$ $1, \ldots d$.

Proof. Let $\eta \in C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$. We calculate

$$
\int_{\mathbb{R}^{d}} \chi^{k}(x, \omega) \partial_{i} \eta(x) d x=\int_{\mathbb{R}^{d}}\left(\sum_{j=1}^{d} \int_{0}^{1} x_{j} U_{j}^{k}\left(\tau_{t x} \omega\right) d t\right) \partial_{i} \eta(x) d x
$$

By changing the order of integration (this can be done since $U_{j}^{k}\left(\tau_{t x} \omega\right) \in L_{\text {loc }}^{1}$ ) and applying the change of variables $y=t x$, we get

$$
\begin{equation*}
\int_{\mathbb{R}^{d}}\left(\sum_{j=1}^{d} \int_{0}^{1} x_{j} U_{j}^{k}\left(\tau_{t x} \omega\right) d t\right) \partial_{i} \eta(x) d x=\int_{0}^{1} \sum_{j=1}^{d} \int_{\mathbb{R}^{d}} \frac{y_{j}}{t^{d+1}} U_{j}^{k}\left(\tau_{y} \omega\right) \partial_{i} \eta\left(\frac{y}{t}\right) d y d t \tag{2.3}
\end{equation*}
$$

Since for $j \neq i$,

$$
\begin{equation*}
\frac{y_{j}}{t^{d+1}} \partial_{i} \eta\left(\frac{y}{t}\right)=\partial_{i}\left(\frac{y_{j}}{t^{d}} \eta\left(\frac{y}{t}\right)\right) \tag{2.4}
\end{equation*}
$$

Lemma 2.3 and (2.4) give

$$
\begin{equation*}
\int_{\mathbb{R}^{d}} \chi^{k}(x, \omega) \partial_{i} \eta(x) d x=\int_{\mathbb{R}^{d}} U_{i}^{k}\left(\tau_{y} \omega\right) \int_{0}^{1} \sum_{j \neq i} \partial_{j}\left(\frac{y_{j}}{t^{d}} \eta\left(\frac{y}{t}\right)\right)+\frac{y_{i}}{t^{d+1}} \partial_{i} \eta\left(\frac{y}{t}\right) d t d y \tag{2.5}
\end{equation*}
$$

Finally, observe that for $y \neq 0$,

$$
\begin{equation*}
\int_{0}^{1} \sum_{j \neq i} \partial_{j}\left(\frac{y_{j}}{t^{d}} \eta\left(\frac{y}{t}\right)\right)+\frac{y_{i}}{t^{d+1}} \partial_{i} \eta\left(\frac{y}{t}\right) d t=-\int_{0}^{1} \frac{1}{t^{d-1}} \frac{d}{d t}\left(\eta\left(\frac{y}{t}\right)\right) d t=-\eta(y) \tag{2.6}
\end{equation*}
$$

Combining (2.5) and (2.6), we have

$$
\int_{\mathbb{R}^{d}} \chi^{k}(x, \omega) \partial_{i} \eta(x) d x=-\int_{\mathbb{R}^{d}} U_{i}^{k}\left(\tau_{x} \omega\right) \eta(x) d x
$$

Hence we have that for each $\eta \in C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$, there exists a $\widehat{\mathbb{P}}$-null set $N_{\eta}$ such that the above equality holds for $\omega \in \widehat{\Omega} \backslash N_{\eta}$. To obtain the desired result, we have to remove the dependency of $\eta$. Because $C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$ is separable with respect to the supremum norm, we can remove this ambiguity considering a countable dense subset $\left\{\eta_{n}\right\}_{n} \subset C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$ and a null set $N=\bigcup_{n} N_{\eta_{n}}$.

Definition 2.5. We say that $u \in \mathcal{F}_{\text {loc }}^{\omega}$ if for any relatively compact set $G \subset \mathbb{R}^{d}$, there exists a function $u_{G} \in \mathcal{F}^{\omega}$ such that $u=u_{G}$ a.e. on $G \cap W^{\prime}(\omega)$.

Proposition 2.6. For $k=1, \ldots d$, the corrector $\chi^{k}(\cdot, \omega) \in \mathcal{F}_{\text {loc }}^{\omega}$ for $\widehat{\mathbb{P}}$-a.e. $\omega \in \Omega$.
Proof. By construction, there exists $\left\{f_{n}\right\}_{n} \subset \mathcal{C}$ such that $\nabla f_{n} \rightarrow U^{k}, n \rightarrow \infty$ in $\widehat{\mathbb{L}}^{2}$. This implies that for any ball $B \subset \mathbb{R}^{d}$,

$$
\begin{aligned}
& \widehat{\mathbb{E}}\left[\int_{B}\left\langle a\left(\tau_{x} \omega\right)\left(\nabla f_{n}\left(\tau_{x} \omega\right)-\nabla \chi(x, \omega)\right), \nabla f_{n}\left(\tau_{x} \omega\right)-\nabla \chi(x, \omega)\right\rangle d x\right] \\
= & \int_{B} \widehat{\mathbb{E}}\left[\left\langle a\left(\tau_{x} \omega\right)\left(\nabla f_{n}\left(\tau_{x} \omega\right)-U^{k}\left(\tau_{x} \omega\right)\right), \nabla f_{n}\left(\tau_{x} \omega\right)-U^{k}\left(\tau_{x} \omega\right)\right\rangle\right] d x \\
\leq & \Lambda|B| \mathbb{E}\left[\left|\nabla f_{n}-U^{k}\right|^{2}\right] \mathbb{P}\left(0 \in W^{\prime}\right)^{-1} \rightarrow 0
\end{aligned}
$$

Observe that $g_{n}(x, \omega)=f_{n}\left(\tau_{x} \omega\right)-f_{n}(\omega)$ belongs to $C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$ and satisfies

$$
g_{n}(x, \omega)=\sum_{i=1}^{d} \int_{0}^{1} x_{i} \partial_{i} f_{n}\left(\tau_{t x} \omega\right) d t
$$

It follows that $g_{n} \rightarrow \chi^{k}$ on $B$ with respect to $\|\cdot\|_{L^{2}(B)}+\mathcal{E}(\cdot, \cdot)$. This implies that $\chi^{k}(\cdot, \omega) \in \mathcal{F}_{\text {loc }}^{\omega} \widehat{\mathbb{P}}$-a.s.

To obtain a martingale decomposition, we introduce a weak notion of the harmonicity.
Definition 2.7. We say that $u \in \mathcal{F}_{\text {loc }}^{\omega}$ is $\mathcal{E}^{\omega}$-harmonic if it satisfies

$$
\mathcal{E}^{\omega}(u, \varphi)=0
$$

for all $\varphi \in C_{c}^{\infty}\left(\overline{W^{\prime}(\omega)}\right)$.
Set $y^{k}(x, \omega)=x_{k}-\chi^{k}(x, \omega)$.
Proposition 2.8. For $k=1, \ldots, d, x \mapsto y^{k}(x, \omega)$ are $\mathcal{E}^{\omega}$-harmonic $\mathbb{P}$-a.s.
Proof. It is enough to show that for any function $\varphi \in C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$,

$$
\begin{equation*}
\int_{W^{\prime}}\left\langle a \nabla y^{k}, \nabla \varphi\right\rangle d x=0 \tag{2.7}
\end{equation*}
$$

since a function in $C_{c}^{\infty}\left(\overline{W^{\prime}(\omega)}\right)$ is embedded in $C_{c}^{\infty}\left(\mathbb{R}^{d}\right)$.
We introduce a subset of $\mathcal{C}$. Define $\mathcal{C}_{0}$ by

$$
\mathcal{C}_{0}=\left\{v \in \mathcal{C} \mid v(\omega)=0 \text { if } 0 \notin W^{\prime}(\omega)\right\} .
$$

Since

$$
L\left(\tau_{x} \omega\right)=\bigcup_{z \in \omega} B(z-x, \rho)=L(\omega)+x,
$$

we have

$$
W^{\prime}\left(\tau_{x} \omega\right)=W^{\prime}(\omega)+x
$$

and hence the condition $0 \notin W^{\prime}\left(\tau_{x} \omega\right)$ is equivalent to $-x \notin W^{\prime}(\omega)$. To prove (2.7), for all $f \in \mathcal{C}_{0}$, using Fubini's theorem we have

$$
\begin{aligned}
& \widehat{\mathbb{E}}\left[f(\omega) \int_{W^{\prime}}\left\langle a\left(\tau_{x} \omega\right) \nabla y^{k}(x, \omega), \nabla \varphi(x)\right\rangle d x\right] \\
& =\widehat{\mathbb{E}}\left[f(\omega) \sum_{i, j} \int_{W^{\prime}} a_{i j}\left(\tau_{x} \omega\right) \partial_{i} y^{k}(x, \omega) \partial_{j} \varphi(x) d x\right] \\
& =\sum_{i, j} \int_{\mathbb{R}^{d}} \partial_{j} \varphi(x) \widehat{\mathbb{E}}\left[f(\omega) a_{i j}\left(\tau_{x} \omega\right) \partial_{i} y^{k}(x, \omega) 1_{\left\{x \in W^{\prime}\right\}}\right] d x .
\end{aligned}
$$

Next using the shift invariance of $\mathbb{P}$, we have

$$
\begin{aligned}
& \sum_{i, j} \int_{\mathbb{R}^{d}} \partial_{j} \varphi(x) \widehat{\mathbb{E}}\left[f(\omega) a_{i j}\left(\tau_{x} \omega\right) \partial_{i} y^{k}(x, \omega) 1_{\left\{x \in W^{\prime}\right\}}\right] d x \\
& =\sum_{i, j} \int_{\mathbb{R}^{d}} \partial_{j} \varphi(x) \mathbb{E}\left[f(\omega) a_{i j}\left(\tau_{x} \omega\right) \partial_{i} y^{k}(x, \omega) 1_{\left\{x \in W^{\prime}\right\}} 1_{\left\{0 \in W^{\prime}\right\}}\right] \mathbb{P}\left(0 \in W^{\prime}\right)^{-1} d x \\
& =\sum_{i, j} \int_{\mathbb{R}^{d}} \partial_{j} \varphi(x) \mathbb{E}\left[f\left(\tau_{-x} \omega\right) a_{i j}(\omega) \partial_{i} y^{k}(0, \omega) 1_{\left\{0 \in W^{\prime}\right\}} 1_{\left\{-x \in W^{\prime}\right\}}\right] \mathbb{P}\left(0 \in W^{\prime}\right)^{-1} d x \\
& =\sum_{i, j} \widehat{\mathbb{E}}\left[a_{i j}(\omega) \partial_{i} y^{k}(0, \omega) \int_{W^{\prime}} f\left(\tau_{-x} \omega\right) 1_{\left\{-x \in W^{\prime}\right\}} \partial_{j} \varphi(x) d x\right] .
\end{aligned}
$$

Since $f \in \mathcal{C}_{0}$, the trace of a function $x \mapsto f\left(\tau_{x} \omega\right)$ to $\partial\left(-W^{\prime}(\omega)\right)$ is identically zero. Hence
by the Gauss-Green formula, we calculate

$$
\begin{aligned}
& \sum_{i, j=1}^{d} \widehat{\mathbb{E}}\left[a_{i j}(\omega) \partial_{i} y^{k}(0, \omega) \int_{-W^{\prime}} f\left(\tau_{x} \omega\right) \partial_{j} \varphi(-x) d x\right] \\
& =-\sum_{i, j=1}^{d} \widehat{\mathbb{E}}\left[a_{i j}(\omega) \partial_{i} y^{k}(0, \omega) \int_{-W^{\prime}} \partial_{j} f\left(\tau_{x} \omega\right) \varphi(-x) d x\right] \\
& =-\sum_{i, j=1}^{d} \widehat{\mathbb{E}}\left[a_{i j}(\omega) \partial_{j} y^{k}(0, \omega) \int_{W^{\prime}} \partial_{i} f\left(\tau_{-x} \omega\right) \varphi(x) d x\right]
\end{aligned}
$$

Again, by the shift invariance, we find that the last line equals to

$$
-\sum_{i, j=1}^{d} \int_{\mathbb{R}^{d}} \varphi(x) \widehat{\mathbb{E}}\left[a_{i j}(\omega) \partial_{i} y^{k}(0, \omega) \partial_{i} f\left(\tau_{-x} \omega\right) 1_{\left\{-x \in W^{\prime}\right\}}\right] d x
$$

Because $f$ belongs to $\mathcal{C}_{0}$, the function $x \mapsto \partial_{i} f\left(\tau_{x} \omega\right)$ vanishes on $\mathbb{R}^{d} \backslash\left(\overline{-W^{\prime}(\omega)}\right)$. Hence we calculate

$$
\begin{aligned}
& =-\sum_{i, j=1}^{d} \int_{\mathbb{R}^{d}} \varphi(x) \widehat{\mathbb{E}}\left[a_{i j}(\omega) \partial_{i} y^{k}(0, \omega) \partial_{i} f\left(\tau_{-x} \omega\right) 1_{\left\{-x \in W^{\prime}\right\}}\right] d x \\
& =-\sum_{i, j=1}^{d} \int_{\mathbb{R}^{d}} \varphi(x) \widehat{\mathbb{E}}\left[a_{i j}(\omega) \partial_{i} y^{k}(0, \omega) \partial_{i} f\left(\tau_{-x} \omega\right)\right] d x \\
& =\int_{\mathbb{R}^{d}} \varphi(x) \widehat{\mathbb{E}}\left[\left\langle a(\omega) \nabla y^{k}(0, \omega), \nabla f\left(\tau_{-x} \omega\right)\right\rangle\right] d x .
\end{aligned}
$$

Since $\tau_{-x} \mathcal{C} \subset \mathcal{C}$ and $\nabla y^{k} \in\left(\widehat{\mathbb{L}}_{\text {pot }}^{2}\right)^{\perp}$, we get

$$
\int_{\mathbb{R}^{d}} \varphi(x) \widehat{\mathbb{E}}\left[\left\langle a \nabla y^{k}(0, \omega), \nabla f\left(\tau_{-x} \omega\right)\right\rangle\right] d x=0
$$

Since $\mathcal{C}_{0} \subset L^{p}(\Omega, \widehat{\mathbb{P}})$ for all $p \geq 1$ densely, it follows that

$$
\int_{W^{\prime}}\left\langle a \nabla y^{k}(x, \omega), \nabla \varphi(x)\right\rangle d x=0, \quad \widehat{\mathbb{P}}-\mathrm{a} . \mathrm{s} .
$$

This ends the proof.
If a function $f$ is harmonic, the process $M_{t}=f\left(X_{t}\right)$ is a martingale from a standard result of Markov processes. To deduce the same result for $\mathcal{E}^{\omega}$-harmonic functions, we will use the following theorem due to Fukushima, Nakao and Takeda [15, Theorem 3.1]. We refer the terminology of Dirichlet forms to [16].

Let $W$ be a Lipschitz domain on $\mathbb{R}^{d}$. Let $\left(Y_{t}, Q_{x}\right)$ be a Hunt process associated with strongly local regular Dirichlet form $(\mathcal{E}, \mathcal{F})$ on $L^{2}(W)$. Fix a point $x_{0} \in W$ and consider the following conditions for the process $\left(Y_{t}, Q_{x}\right)$ and for a function $u$ :
(1) The transition function $p_{t}$ of $X_{t}$ satisfies $p_{t}\left(x_{0}, A\right)=0$ for any $t$ if $\operatorname{Cap}(A)=0$.
(2) The function $u$ belongs to $\mathcal{F}_{\text {loc }}$. Moreover, the function $u$ is continuous and $\mathcal{E}$ harmonic.
(3) The energy measure $\nu_{\langle u\rangle}$ of $u$ is absolutely continuous with respect to the Lebesgue measure on $W$ and the density function $f$ satisfies

$$
E_{x_{0}}\left[\int_{0}^{t} f\left(X_{s}\right) d x\right]<\infty, t>0
$$

Lemma 2.9. ([15, Theorem 3.1]) Assume that the above conditions hold. Then the additive functional $M_{t}=u\left(X_{t}\right)-u\left(X_{0}\right)$ is a $P_{x_{0}}$-square integrable martingale with

$$
\begin{equation*}
\langle M\rangle_{t}=\int_{0}^{t} f\left(X_{s}\right) d s, \quad t>0, \quad P_{x_{0}-\text { a.s. }} \tag{2.8}
\end{equation*}
$$

Corollary 2.10. The process $y\left(X_{t}, \omega\right)$ is a martingale with quadratic variation

$$
\left\langle y^{k}(X ., \omega), y^{\ell}(X ., \omega)\right\rangle_{t}=\int_{0}^{t} \sum_{i, j=1}^{d} a_{i j}\left(\tau_{x} \omega\right) \partial_{i} y^{k}\left(X_{t}, \omega\right) \partial_{j} y^{\ell}\left(X_{t}, \omega\right) d s,(k, \ell=1, \ldots, d)
$$

Proof. To apply Lemma 2.9 for $y^{k}$, we need to check the assumptions of Lemma 2.9.
Since the reflecting diffusion $\left\{X_{t}\right\}_{t}$ has the transition density (say $q_{t}(x, y)$ ), we have $p_{t}\left(x_{0}, A\right)=\int_{A} q_{t}\left(x_{0}, y\right) d y=0$ if Cap $(A)=0$. Hence assumption (1) of Lemma 2.9 holds.

By Proposition 2.8, for $\widehat{\mathbb{P}}$-a.s. $\omega$, the function $x \mapsto y^{k}(x, \omega)$ satisfies a solution of the equation

$$
\int_{W}\left\langle a\left(\tau_{x} \omega\right) \nabla y^{k}(x, \omega), \nabla \varphi(x)\right\rangle d x=0, \quad \varphi \in C_{c}^{\infty}(W)
$$

Thus, assumption (2) follows from classical results in elliptic partial differential equations (see [19, Section 9]). We have

$$
\begin{aligned}
& 2 \mathcal{E}^{\omega}\left(y^{k} v, y^{k}\right)-\mathcal{E}^{\omega}\left(\left(y^{k}\right)^{2}, v\right) \\
& =2 \int_{W^{\prime}} v\left\langle a\left(\tau_{x} \omega\right) \nabla y^{k}, \nabla y^{k}\right\rangle d x+2 \int_{W^{\prime}} y^{k}\left\langle a\left(\tau_{x} \omega\right) \nabla v, \nabla y^{k}\right\rangle d x-2 \int_{W^{\prime}} y^{k}\left\langle a\left(\tau_{x} \omega\right) \nabla y^{k}, \nabla v\right\rangle d x \\
& =\int_{W^{\prime}} 2\left\langle a\left(\tau_{x} \omega\right) \nabla y^{k}, \nabla y^{k}\right\rangle v d x
\end{aligned}
$$

for all $v \in C_{c}^{\infty}(W)$. Since the function $x \mapsto y^{k}(x, \omega)$ is weakly differentiable, the energy measure $\nu_{\left\langle y^{k}\right\rangle}$ is given by

$$
\nu_{\left\langle y^{k}\right\rangle}(d x)=2\left\langle a\left(\tau_{x} \omega\right) \nabla y^{k}(x), \nabla y^{k}(x)\right\rangle d x
$$

and absolutely continuous with respect to the Lebesgue measure. Moreover, its density $f$ is written by

$$
f(x)=2\left\langle a\left(\tau_{x} \omega\right) \nabla y^{k}(x), \nabla y^{k}(x)\right\rangle
$$

Next using the stationarity of the environment process $\left\{\tau_{X_{t}^{\omega}}\right\}_{t}$ under $\widehat{\mathbb{P}}$ (see [26]), by
using Fubini's Theorem and the translation invariance, we compute

$$
\widehat{\mathbb{E}}\left[E_{0}^{\omega}\left[\int_{0}^{t} f\left(X_{s}^{\omega}\right) d s\right]\right]=\int_{0}^{t} \widehat{\mathbb{E}}\left[E_{0}^{\omega}\left[f\left(X_{s}^{\omega}\right)\right]\right] d s=\int_{0}^{t} \widehat{\mathbb{E}}\left[E_{0}^{\omega}\left[f\left(X_{0}^{\omega}\right)\right]\right] d s=t \widehat{\mathbb{E}}\left[E_{0}^{\omega}[f(0)]\right]<\infty
$$

Therefore, assumption (3) is satisfied.

## 3. Maximal inequality

In this section, we consider deterministic settings. Let $W \subset \mathbb{R}^{d}$ be an unbounded Lipschitz domain containing the origin. We denote the connected component of $B(0, R) \cap$ $W$ containing the origin by $W_{R}$. We assume that there exists positive constants $C_{V}$ and $R_{V}$ such that

$$
\begin{equation*}
\left|W_{R}\right| \geq C_{V} R^{d} \tag{3.1}
\end{equation*}
$$

for $R \geq R_{V}$. We further assume that there exist $\theta \in(0,1)$ and $c_{H}>0$ such that

$$
C_{\mathrm{IL}}:=\inf \left\{\frac{\mathcal{H}_{d-1}(W \cap \partial O)}{|W \cap O|^{\frac{d-1}{d}}} \left\lvert\, \begin{array}{l}
O \subset B(0, R) \text { is connected open, }  \tag{3.2}\\
R \geq R_{I},|O| \geq R^{\theta}
\end{array}\right.\right\}>0
$$

for some $R_{I}>0$ and

$$
C_{\mathrm{IS}}:=\inf \left\{\frac{\mathcal{H}_{d-1}(W \cap \partial O)}{|W \cap O|^{\frac{d-1}{d}}} \left\lvert\, \begin{array}{l}
O \text { is bounded open, }  \tag{3.3}\\
\mathcal{H}_{d-1}(W \cap \partial O)<c_{H}
\end{array}\right.\right\}>0
$$

hold.
Set $\zeta=\frac{1-\theta}{1-\frac{\theta}{d}}$. For a bounded open subset, the following weak isoperimetric inequality holds.

Lemma 3.1. There exists positive constant $C_{I}$ such that

$$
\begin{equation*}
\frac{\mathcal{H}_{d-1}(W \cap \partial O)}{|O|^{\frac{d-\zeta}{d}}} \geq \frac{C_{I}}{R^{1-\zeta}} \tag{3.4}
\end{equation*}
$$

holds for open subset $O \subset W(0, R)$ and $R \geq R_{I}$.
Proof. The general idea of the proof goes back to [9, Lemma, 3.3]. First we consider a good and connected open subset $O$. When $O$ satisfies $|O| \geq R^{\theta}$ or $\mathcal{H}_{d-1}(W \cap \partial O)<c_{H}$, from (3.2) and (3.3), we have

$$
\begin{equation*}
\frac{\mathcal{H}_{d-1}(W \cap \partial O)}{|O|^{\frac{d-\zeta}{d}}}=\frac{\mathcal{H}_{d-1}(W \cap \partial O)}{|O|^{\frac{d-1}{d}}|O|^{\frac{1-\zeta}{d}}} \geq \frac{C_{\mathrm{IL}} \wedge C_{\mathrm{IS}}}{|O|^{\frac{1-\zeta}{d}}} \geq \frac{C_{\mathrm{IL}} \wedge C_{\mathrm{IS}}}{|B(0, R)|^{\frac{1-\zeta}{d}}} \geq \frac{C}{R^{1-\zeta}} \tag{3.5}
\end{equation*}
$$

Next we consider the case where $|O|<R^{\theta}$ and $\mathcal{H}_{d-1}(W \cap \partial O) \geq c_{H}$. Using the relation $\theta \cdot \frac{d-\zeta}{d}=1-\zeta$, we estimate

$$
\begin{equation*}
\frac{\mathcal{H}_{d-1}(W \cap \partial O)}{|O|^{\frac{d-\zeta}{d}}} \geq \frac{c_{H}}{|O|^{\frac{d-\zeta}{d}}} \geq \frac{c_{H}}{R^{\theta \cdot \frac{d-\zeta}{d}}}=\frac{c_{H}}{R^{1-\zeta}} \tag{3.6}
\end{equation*}
$$

Combining (3.5) and (3.6), we get (3.4) when $O$ is connected. For a general open subset $O$, it follows from the fact that $\left(a_{1}+a_{2}\right) /\left(b_{1}+b_{2}\right)^{\beta} \geq a_{1} / b_{1}^{\beta}+a_{2} / b_{2}^{\beta} \geq M$ for $\beta \in(0,1)$ and positive numbers $a_{1}, a_{2}, b_{1}, b_{2}, M$ with $a_{1} / b_{1}^{\beta} \geq M$ and $a_{2} / b_{2}^{\beta} \geq M$.

Let $E \subset W$ be a bounded set. For a function $u: E \rightarrow \mathbb{R}, \alpha \geq 1$ we denote

$$
\|u\|_{L^{\alpha}(E)}:=\left(\int_{E}|u(x)|^{\alpha} d x\right)^{\frac{1}{\alpha}},\|u\|_{E, \alpha}:=\left(\frac{1}{|E|} \int_{E}|u(x)|^{\alpha} d x\right)^{\frac{1}{\alpha}}
$$

Next we show a Sobolev type inequality. To do this, we need to consider functions whose trace vanishes locally. For $R>0$, set

$$
\mathcal{C}_{R}=\left\{u \in C_{c}^{\infty}\left(\overline{W_{R}}\right) \mid u=0 \text { on } \partial W_{R}-\partial W\right\} .
$$

Let $a: \mathbb{R}^{d} \rightarrow \mathbb{R}^{d \times d}$ be a positive-definite symmetric matrix. Define the bilinear form $\mathcal{E}$ on $L^{2}(W, d x)$ by

$$
\begin{equation*}
\mathcal{E}(u, v)=\int_{W}\langle a \nabla u, \nabla v\rangle d x, \quad u, v \in L^{2}(W, d x) \cap C^{\infty}(W) \tag{3.7}
\end{equation*}
$$

Assumption 5. (1) there exist positive measurable functions $\lambda, \Lambda$ such that for almost all $x, \xi \in \mathbb{R}^{d}$,

$$
\lambda(x)|\xi|^{2} \leq\langle a(x) \xi, \xi\rangle \leq \Lambda(x)|\xi|^{2}
$$

(2) there exist $p, q \in[1, \infty]$ satisfying $1 / p+1 / q<2 \zeta / d$ such that for almost all $x$,

$$
\limsup _{R \rightarrow \infty} \frac{1}{\left|W_{R}\right|} \int_{W_{R}}\left(\Lambda^{p}+\lambda^{-q}\right) d x<\infty
$$

We remark that the relation between $p$ and $q$ is more restricted than Chiarini and Deuschel [9]. This is because of the boundary effect. Let $\mathcal{F}_{R}$ be the closure of $\mathcal{C}_{R}$ with respect to $\|\cdot\|_{L^{2}\left(\mathcal{C}_{R}\right)}+\mathcal{E}(\cdot, \cdot)$. Then if the Dirichlet form $\left(\mathcal{E}, \mathcal{F}_{R}\right)$ is regular, the associate diffusion is absorbed at $\partial W_{R}-\partial W$ and has a reflection at $\partial W_{R} \cap \partial W$. We note that the domain $\mathcal{C}_{R}$ is larger than the Sobolev space with zero boundary condition $H_{0}^{1}\left(W_{R}\right)$. Hence, it is not obvious whether the Sobolev inequality holds. The key property is the weak isoperimetric inequality (3.4).

Proposition 3.2. Let $R \geq R_{I} \vee R_{V}$. Then we have

$$
\begin{equation*}
\|u\|_{L^{\frac{d}{d-\zeta}}\left(W_{R}\right)} \leq C_{S}^{-1}\left|W_{R}\right|^{\frac{1-\zeta}{d}}\|\nabla u\|_{L^{1}\left(W_{R}\right)} \tag{3.8}
\end{equation*}
$$

for $u \in \mathcal{F}_{R}$, where $C_{S}=C_{I} C_{V}^{\frac{1-\zeta}{d}}$.
Proof. We first show (3.8) for non-negative $u \in \mathcal{C}_{R}$. Let $v: \mathbb{R}^{d} \rightarrow \mathbb{R}$ be a zero-extension of $u$. That is, $v=u$ in $\operatorname{supp} u$ and $v=0$ outside the $\operatorname{supp} u$. Then by the coarea formula
(see [11, Theorem 3.10]), we have

$$
\int_{W_{R}}|\nabla v| d x=\int_{-\infty}^{\infty} \mathcal{H}_{d-1}\left(W_{R} \cap v^{-1}(\{t\})\right) d t
$$

Now, since $v$ is non-negative, we have $v^{-1}(\{t\})=\varnothing$ for $t<0$. Moreover, $u=v$ on $W_{R}$. Thus, we have

$$
\begin{equation*}
\int_{W_{R}}|\nabla u| d x=\int_{0}^{\infty} \mathcal{H}_{d-1}\left(W_{R} \cap \partial\left\{x \in W_{R} \mid u(x)>t\right\}\right) d t \tag{3.9}
\end{equation*}
$$

Since $u$ is continuous and $\min u=0$, we have $W_{R} \cap\left\{x \in W_{R} \mid u(x)=t\right\}=W_{R} \cap \partial\{x \in$ $\left.W_{R} \mid u(x)>t\right\}$ for $t>0$. Since $u=0$ on $\partial W_{R}-\partial W$, we also have $W_{R} \cap \partial\left\{x \in W_{R} \mid\right.$ $u(x)>t\}=W \cap \partial\left\{x \in W_{R} \mid u(x)>t\right\}$. Combining these and (3.9), we obtain

$$
\begin{equation*}
\int_{W_{R}}|\nabla u| d x=\int_{0}^{\infty} \mathcal{H}_{d-1}\left(W \cap \partial\left\{x \in W_{R} \mid u(x)>t\right\}\right) d t \tag{3.10}
\end{equation*}
$$

Set $U_{t}=\left\{x \in W_{R} \mid u(x)>t\right\}$. From (3.4), we estimate

$$
\begin{align*}
& \int_{0}^{\infty} \mathcal{H}_{d-1}\left(W \cap \partial\left\{x \in W_{R} \mid u(x)>t\right\}\right) d t \\
& \geq \frac{C_{I}}{R^{1-\zeta}} \int_{0}^{\infty}\left|W \cap\left\{x \in W_{R} \mid u(x)>t\right\}\right|^{\frac{d-\zeta}{d}} d t \\
& =\frac{C_{I}}{R^{1-\zeta}} \int_{0}^{\infty}\left|\left\{x \in W_{R} \mid u(x)>t\right\}\right|^{\frac{d-\zeta}{d}} d t \\
& =\frac{C_{I}}{R^{1-\zeta}} \int_{0}^{\infty}\left|U_{t}\right|^{\frac{d-\zeta}{d}} d t \\
& =\frac{C_{I}}{R^{1-\zeta}} \int_{0}^{\infty}\left\|1_{U_{t}}\right\|_{L^{\frac{d}{d-\zeta}}\left(W_{R}\right)} d t . \tag{3.11}
\end{align*}
$$

Using (3.1), we have $1 / R^{1-\zeta} \geq\left(C_{V} /\left|W_{R}\right|\right)^{\frac{1-\zeta}{d}}$. Combining this with (3.10) and (3.11), we obtain

$$
\begin{equation*}
\int_{W_{R}}|\nabla u| d x \geq \frac{C_{S}}{\left|W_{R}\right|^{\frac{1-\zeta}{d}}} \int_{0}^{\infty}\left\|1_{U_{t}}\right\|_{L^{\frac{d}{d-\zeta}\left(W_{R}\right)}} d t \tag{3.12}
\end{equation*}
$$

Now let $r$ be positive number satisfying $(d /(d-\zeta))^{-1}+1 / r=1$. Take $g \in L^{r}$ such that $g \geq 0$ and $\|g\|_{L^{r}\left(W_{R}\right)}=1$. Then, by the Hölder inequality,

$$
\int_{0}^{\infty}\left\|1_{U_{t}}\right\|_{L^{\frac{d}{d-1}}\left(W_{R}\right)} d t \geq \int_{0}^{\infty}\left\|g 1_{U_{t}}\right\|_{L^{1}\left(W_{R}\right)} d t=\int_{W_{R}} g(x) \int_{0}^{\infty} 1_{U_{t}}(x) d t d x=\|g u\|_{L^{1}\left(W_{R}\right)}
$$

Since $g$ is arbitrary, putting this into (3.12), we obtain

$$
\|u\|_{L^{\frac{d}{d-\zeta}\left(W_{R}\right)}}=\sup _{g \in L^{r}\left(W_{R}\right), g \geq 0,\|g\|_{L^{r}\left(W_{R}\right)}=1}\|g u\|_{L^{1}\left(W_{R}\right)} \leq C_{S}^{-1}\left|W_{R}\right|^{\frac{1-\zeta}{d}}\|\nabla u\|_{L^{1}\left(W_{R}\right)}
$$

hence (3.8) holds. For general $u \in \mathcal{F}_{R}$, we can show the same bound by approximation.

Set $p^{*}=\frac{p}{p-1}, \rho=\frac{2 q d}{q(d-2 \zeta)+d}$, and $q^{\#}=\frac{2 q}{q+1}$.
Proposition 3.3. Let $R \geq R_{I} \vee R_{V}$. Let $p, q$ be positive numbers as in (2) of Assumption 5. Then we have

$$
\begin{equation*}
\|u\|_{L^{\rho}\left(W_{R}\right)} \leq C_{\mathrm{sob}}\left|W_{R}\right|^{\frac{1-\zeta}{d}}\|\nabla u\|_{L^{q} \#}\left(W_{R}\right) \tag{3.13}
\end{equation*}
$$

for $u \in \mathcal{F}_{R}$.
Proof. First observe that $\rho=\frac{d q^{\#}}{d-q^{\#} \zeta}$. Let $r$ be a real number satisfying $1 / q^{\#}+1 / r=1$.
Set $\bar{d}=q^{\#}(d-\zeta) /\left(d-q^{\#} \zeta\right)$. Using a relation $\frac{d q^{\#}}{d-q^{\#} \zeta}=\frac{q^{\#}(d-\zeta)}{d-q^{\#} \zeta} \frac{d}{d-\zeta}=\bar{d} \cdot \frac{d}{d-\zeta}$ and Proposition 3.2, we have

$$
\begin{align*}
& \left(\int_{E}|u|^{\frac{d q}{d-q^{\#}}} d x\right)^{\frac{d-1}{d}}=\left\|u^{\bar{d}}\right\|_{L^{\frac{d}{d-\zeta}}\left(W_{R}\right)} \\
& \leq C_{S}^{-1}\left|W_{R}\right|^{\frac{1-\zeta}{d}}\left\|\nabla u^{\bar{d}}\right\|_{L^{1}\left(W_{R}\right)}=\bar{d} C_{S}^{-1}\left|W_{R}\right|^{\frac{1-\zeta}{d}}\left\|u^{\bar{d}-1} \nabla u\right\|_{L^{1}\left(W_{R}\right)} \tag{3.14}
\end{align*}
$$

Using the Hölder inequality with a pair $\left(r, q^{\#}\right)$ and a relation $(\bar{d}-1) r=\frac{d q^{\#}}{d-q^{\#} \zeta}$, we have

$$
\begin{aligned}
& \bar{d} C_{S}^{-1}\left|W_{R}\right|^{\frac{1-\zeta}{d}}\left\|u^{\bar{d}-1} \nabla u\right\|_{L^{1}\left(W_{R}\right)} \leq \bar{d} C_{S}^{-1}\left|W_{R}\right|^{\frac{1-\zeta}{d}}\left\|u^{\bar{d}-1}\right\|_{L^{r}\left(W_{R}\right)}\|\nabla u\|_{L^{q} \#}\left(W_{R}\right) \\
& =\bar{d} C_{S}^{-1}\left|W_{R}\right|^{\frac{1-\zeta}{d}}\|\nabla u\|_{L^{q \#}\left(W_{R}\right)}\left(\int_{W_{R}}|u|^{\frac{d q^{\#}}{d-q^{\#}}} d x\right)^{\frac{q^{\#-1}}{q^{\#}}}
\end{aligned}
$$

Inserting this into (3.14), we get

$$
\left(\int_{W_{R}}|u|^{\frac{d q^{\#}}{d-q^{\#} \zeta}} d x\right)^{\frac{d-1}{d}} \leq \bar{d} C_{S}^{-1}\left|W_{R}\right|^{\frac{1-\zeta}{d}}\|\nabla u\|_{L^{q} \#}\left(W_{R}\right)\left(\int_{W_{R}}|u|^{\frac{d q^{\#}}{d-q^{\#} \zeta}} d x\right)^{\frac{q^{\#}-1}{q^{\#}}}
$$

Dividing both side by $\left(\int|u|^{\frac{d q^{\#}}{d-q \#}} d x\right)^{\frac{q^{\#}-1}{q^{\#}}}$, we get the desired result since $(d-\zeta) / d-$ $\left(q^{\#}-1\right) / q^{\#}=\left(d-q^{\#} \zeta\right) /\left(d q^{\#}\right)$.

Thanks to Proposition 3.3, we can prove the following inequality.
Lemma 3.4. Let $R \geq R_{I} \vee R_{V}$. Then for all $u \in \mathcal{F}_{R}$

$$
\begin{equation*}
\|u\|_{L^{\rho}\left(W_{R}\right)}^{2} \leq C_{\mathrm{sob}}\left|W_{R}\right|^{\frac{1-\zeta}{d}}\left\|\lambda^{-1}\right\|_{L^{q}\left(W_{R}\right)} \mathcal{E}(u, u) \tag{3.15}
\end{equation*}
$$

Proof. By (3.13), we have

$$
\|u\|_{L^{\rho}\left(W_{R}\right)} \leq C_{\mathrm{sob}}\left|W_{R}\right|^{\frac{1-\zeta}{d}}\|\nabla u\|_{L^{q^{\#}}\left(W_{R}\right)} .
$$

By the Hölder inequality and (1) of Assumption 5, the right hand side is estimated by

$$
\|\nabla u\|_{L^{q} \#}\left(W_{R}\right)=\left(\int_{W_{R}}|\nabla u|^{\frac{2 q}{q+1}} \lambda^{\frac{q}{q+1}} \cdot \lambda^{-\frac{q}{q+1}} d x\right)^{\frac{q+1}{2 q}}
$$

$$
\begin{aligned}
& \leq\left(\int_{W_{R}}\left(|\nabla u|^{2} \lambda\right)^{\frac{q}{q+1} \cdot \frac{q+1}{q}} d x\right)^{\frac{q+1}{2 q} \cdot \frac{q}{q+1}}\left\|\lambda^{-1}\right\|_{L^{q}\left(W_{R}\right)} \\
& \leq \mathcal{E}(u, u)^{\frac{1}{2}}\left\|\lambda^{-1}\right\|_{L^{q}\left(W_{R}\right)} .
\end{aligned}
$$

Therefore, we get the desired result.
By cutoff on $W_{R}$ we mean a function $\eta \in \mathcal{C}_{R}$ satisfying $0 \leq \eta \leq 1$.
Proposition 3.5. Let $R \geq R_{I} \vee R_{V}$. Let $\eta$ be a cutoff on $W_{R}$. Then there exists a constant $C>0$, depending only on the dimension $d \geq 2$, such that for all $u \in \mathcal{F}_{R}$

$$
\begin{equation*}
\|\eta u\|_{L^{\rho}\left(W_{R}\right)}^{2} \leq 2 C\left|W_{R}\right|^{\frac{1-\zeta}{d}}\left\|\lambda^{-1}\right\|_{L^{q}\left(W_{R}\right)}\left(\mathcal{E}_{\eta}(u, u)+\|\eta\|_{\infty}^{2}\left\|u \Lambda^{\frac{1}{2}}\right\|_{L^{2}\left(W_{R}\right)}^{2}\right) \tag{3.16}
\end{equation*}
$$

where we denote $\mathcal{E}_{\eta}(u, u)=\int_{W}\langle a \nabla u, \nabla u\rangle \eta^{2} d x$.
Proof. By Lemma 3.4, we have

$$
\|\eta u\|_{\rho}^{2} \leq C_{\mathrm{sob}}\left|W_{R}\right|^{\frac{1-\zeta}{d}}\left\|\lambda^{-1}\right\|_{L^{q}\left(W_{R}\right)} \mathcal{E}(\eta u, \eta u)
$$

Because $\langle a(\eta \nabla u+u \nabla \eta),(\eta \nabla u+u \nabla \eta)\rangle \leq 2\left(\eta^{2}\langle a \nabla u, \nabla u\rangle+u^{2}\langle a \nabla \eta, \nabla \eta\rangle\right)$, we estimate

$$
\begin{aligned}
\mathcal{E}(\eta u, \eta u) & =\int_{W}\langle a \nabla(\eta u), \nabla(\eta u)\rangle d x \\
& =\int_{W}\langle a(\eta \nabla u+u \nabla \eta),(\eta \nabla u+u \nabla \eta)\rangle d x \\
& \leq \int_{W} 2\left(\eta^{2}\langle a \nabla u, \nabla u\rangle+u^{2}\langle a \nabla \eta, \nabla \eta\rangle\right) d x \\
& \leq 2 \int_{W}\langle a \nabla u, \nabla u\rangle \eta^{2} d x+2 \int_{W}\|\nabla \eta\|_{\infty}^{2} u^{2} \Lambda d x \\
& =2 \mathcal{E}_{\eta}(u, u)+2\|\eta\|_{\infty}^{2}\left\|u \Lambda^{\frac{1}{2}}\right\|_{L^{2}\left(W_{R}\right)}^{2},
\end{aligned}
$$

which leads to the conclusion.
Let $\mathcal{F}$ be the closure of $C_{c}^{\infty}(\bar{W})$ with respect to $\|\cdot\|_{L^{2}(W)}+\mathcal{E}(\cdot, \cdot)$. We say that a function $u$ belongs to $\mathcal{F}_{\text {loc }}$ if for all $R>0$ there exists $u_{R} \in \mathcal{F}$ such that $u=u_{R}$ on $W_{R}$. Let $f: W \rightarrow \mathbb{R}$ be a function with essentially bounded derivatives. Consider the following equation:

$$
\begin{equation*}
\mathcal{E}(u, \varphi)=-\int_{W}\langle a \nabla f, \nabla \varphi\rangle d x \tag{3.17}
\end{equation*}
$$

We say that $u \in \mathcal{F}_{\text {loc }}$ is a solution of the equation (3.17) if it holds for all $\varphi \in C_{c}^{\infty}(\bar{W})$. We say that $u$ is a subsolution of the equation (3.17) if the equation (3.17) holds with $\leq$ for all $\varphi \in C_{c}^{\infty}(\bar{W})$. We also say that $u$ is a solution of the equation in $W_{R}$ if the equation (3.17) holds for $\varphi \in \mathcal{F}_{R}$.

Proposition 3.6. Let $R \geq R_{I} \vee R_{V}$ and $u \in \mathcal{F}_{\text {loc }}$ be a subsolution of equation of (3.17) in $W_{R}$. Let $\eta \in C_{c}^{\infty}\left(W_{R}\right)$ be a cutoff. Then there exists a constant $C_{1}>0$ such
that for all $\alpha \geq 1$,

$$
\begin{align*}
\|\eta u\|_{W_{R}, \alpha \rho}^{2 \alpha} \leq & \alpha^{2} C_{1}\left\|\lambda^{-1}\right\|_{W_{R}, q}\|\Lambda\|_{W_{R}, p}\left|W_{R}\right|^{\frac{2}{d}}  \tag{3.18}\\
& \times\left(\|\nabla \eta\|_{\infty}^{2}\left\|u^{+}\right\|_{W_{R}, 2 \alpha p^{*}}^{2 \alpha}+\|\nabla f\|_{\infty}^{2}\left\|u^{+}\right\|_{W_{R}, 2 \alpha p^{*}}\right)
\end{align*}
$$

Proof. Similarly to [9, Proposition 2.4], we can show that

$$
\begin{aligned}
\left\|\eta\left(u^{+}+\varepsilon\right)^{\alpha}\right\|_{\rho}^{2} & \leq 2 C\left|W_{R}\right|^{\frac{1-\zeta}{d}}\left\|\lambda^{-1}\right\|_{L^{q}\left(W_{R}\right)}\left[\left(\alpha^{2}+1\right)\left\|\left(u^{+}+\varepsilon\right)^{2 \alpha} \Lambda\right\|_{L^{1}\left(W_{R}\right)}\|\nabla \eta\|_{L^{\infty}\left(W_{R}\right)}^{2}\right. \\
& +\|\nabla f\|_{L^{\infty}\left(W_{R}\right)}^{2} \alpha^{2}\left\|\left(u^{+}+\varepsilon\right)^{2 \alpha-2} \Lambda\right\|_{L^{1}\left(W_{R}\right)} \\
& \left.+\frac{\alpha^{2}}{2 \alpha-1}\|\nabla \eta\|_{L^{\infty}\left(W_{R}\right)}\|\nabla f\|_{L^{\infty}\left(W_{R}\right)}\left\|u^{2 \alpha-1} \Lambda\right\|_{L^{1}\left(W_{R}\right)}\right] .
\end{aligned}
$$

(The only difference is using (3.16) instead of [9, Proposiotion 2.3].). Taking the limit as $\varepsilon \rightarrow 0$ and using the Hölder inequality with $1 / p+1 / p^{*}=1$, we have

$$
\begin{aligned}
\left\|\eta\left(u^{+}\right)^{\alpha}\right\|_{\rho}^{2} & \leq 2 C\left|W_{R}\right|^{\frac{1-\zeta}{d}}\left\|\lambda^{-1}\right\|_{L^{q}\left(W_{R}\right)}\left[\left(\alpha^{2}+1\right)\left\|\left(u^{+}\right)^{2 \alpha}\right\|_{L^{p^{*}}\left(W_{R}\right)}\|\Lambda\|_{L^{p}\left(W_{R}\right)}\|\nabla \eta\|_{L^{\infty}\left(W_{R}\right)}^{2}\right. \\
& +\|\nabla f\|_{L^{\infty}\left(W_{R}\right)}^{2} \alpha^{2}\left\|\left(u^{+}\right)^{2 \alpha-2} \Lambda\right\|_{L^{p^{*}}\left(W_{R}\right)}\|\Lambda\|_{L^{p}\left(W_{R}\right)} \\
& \left.+\frac{\alpha^{2}}{2 \alpha-1}\|\nabla \eta\|_{L^{\infty}\left(W_{R}\right)}\|\nabla f\|_{L^{\infty}\left(W_{R}\right)}\left\|u^{2 \alpha-1}\right\|_{L^{p^{*}\left(W_{R}\right)}}\|\Lambda\|_{L^{p}\left(W_{R}\right)}\right] .
\end{aligned}
$$

Averaging over $W_{R}$ and using the relation $2 / \rho=1 / p^{*}+1 / p+1 / q-(2 \zeta) / d$ and $(1+\zeta) / d \leq 2 / d$, we get

$$
\begin{aligned}
\left\|\eta\left(u^{+}\right)^{\alpha}\right\|_{W_{R}, \rho}^{2} & \leq 2 C\left\|\lambda^{-1}\right\|_{W_{R}, q}\|\Lambda\|_{W_{R}, p}\left|W_{R}\right|^{\frac{1+\zeta}{d}}\left[\left(\alpha^{2}+1\right)\left\|\left(u^{+}\right)^{2 \alpha}\right\|_{W_{R}, p^{*}}\|\nabla \eta\|_{\infty}^{2}\right. \\
& \left.+\alpha^{2}\|\nabla f\|_{\infty}^{2}\left\|\left(u^{+}\right)^{2 \alpha-2}\right\|_{W_{R}, p^{*}}+\frac{\alpha^{2}}{2 \alpha-1}\|\nabla \eta\|_{\infty}\|\nabla f\|_{\infty}\left\|\left(u^{+}\right)^{2 \alpha-1}\right\|_{W_{R}, p^{*}}\right] \\
& \leq 2 C\left\|\lambda^{-1}\right\|_{W_{R}, q}\|\Lambda\|_{W_{R}, p}\left|W_{R}\right|^{\frac{2}{d}}\left[\left(\alpha^{2}+1\right)\left\|\left(u^{+}\right)^{2 \alpha}\right\|_{W_{R}, p^{*}}\|\nabla \eta\|_{\infty}^{2}\right. \\
& \left.+\alpha^{2}\|\nabla f\|_{\infty}^{2}\left\|\left(u^{+}\right)^{2 \alpha-2}\right\|_{W_{R}, p^{*}}+\frac{\alpha^{2}}{2 \alpha-1}\|\nabla \eta\|_{\infty}\|\nabla f\|_{\infty}\left\|\left(u^{+}\right)^{2 \alpha-1}\right\|_{W_{R}, p^{*}}\right] .
\end{aligned}
$$

By Jensen's inequality we have

$$
\left\|u^{+}\right\|_{W_{R},(2 \alpha-2) p^{*}} \leq\left\|u^{+}\right\|_{W_{R}, 2 \alpha p^{*}}, \quad\left\|u^{+}\right\|_{W_{R},(2 \alpha-1) p^{*}} \leq\left\|u^{+}\right\|_{W_{R}, 2 \alpha p^{*}},
$$

therefore we can rewrite and get

$$
\begin{aligned}
\left\|\eta u^{+}\right\|_{W_{R}, \rho}^{2} & \leq 2 C\left|W_{R}\right|^{\frac{2}{d}}\left[\left(\alpha^{2}+1\right)\left\|1_{W_{R}} u^{+}\right\|_{W_{R}, 2 \alpha p^{*}}^{2 \alpha}\|\nabla \eta\|_{\infty}^{2}\right. \\
& \left.+\alpha^{2}\|\nabla f\|_{\infty}^{2}\left\|1_{W_{R}} u^{+}\right\|_{W_{R}, 2 \alpha p^{*}}^{2 \alpha-2}+\frac{\alpha^{2}}{2 \alpha-1}\|\nabla \eta\|_{\infty}\|\nabla f\|_{\infty}\left\|u^{+}\right\|_{W_{R}, 2 \alpha p^{*}}^{2 \alpha-1}\right]
\end{aligned}
$$

Absorbing the mixed product in the two squares, we obtain the desired result.
We can prove the following inequality as in [9, Corollary 2.1].

Corollary 3.7. Let $R \geq R_{I} \vee R_{V}$ and $u \in \mathcal{F}_{\text {loc }}$ be a solution of equation of (3.17) in $W_{R}$. Let $\eta \in C_{c}^{\infty}\left(W_{R}\right)$ be a cutoff function. Then there exists a constant $C_{1}>0$ such that for all $\alpha \geq 1$,

$$
\begin{align*}
\|\eta u\|_{W_{R}, \alpha \rho}^{2 \alpha} \leq & \alpha^{2} C_{1}\left\|\lambda^{-1}\right\|_{W_{R}, q}\|\Lambda\|_{W_{R}, p}\left|W_{R}\right|^{\frac{2}{d}}  \tag{3.19}\\
& \times\left(\|\nabla \eta\|_{\infty}^{2}\|u\|_{W_{R}, 2 \alpha p^{*}}^{2 \alpha}+\|\nabla f\|_{\infty}^{2}\|u\|_{W_{R}, 2 \alpha p^{*}}\right) .
\end{align*}
$$

The general idea of the proof of the following proposition is similar to that of [9, Corollary 2.2 ] but we need to use (3.1).

Proposition 3.8. Take $R \geq R_{I} \vee R_{V}$ so that (3.1) and (3.4) hold. We write $W(R)$ for $W_{R}$. Suppose that $u$ is a solution of (3.17) in $W(R)$, and assume that $|\nabla f| \leq c_{f} / R$. Then there exist $\kappa \in(1, \infty), \gamma \in(0,1]$ and $C_{2}=C_{2}\left(c_{f}\right)>0$ such that

$$
\begin{equation*}
\|u\|_{W\left(\sigma^{\prime} R\right), \infty} \leq C_{2}\left(\frac{1 \vee\left\|\lambda^{-1}\right\|_{W(R), q}\|\Lambda\|_{W(R), p}}{\left(\sigma-\sigma^{\prime}\right)^{2}}\right)^{\kappa}\|u\|_{W(\sigma R), \rho}^{\gamma} \vee\|u\|_{W(\sigma R), \rho} \tag{3.20}
\end{equation*}
$$

for any fixed $1 / 2 \leq \sigma^{\prime}<\sigma \leq 1$.
Proof. Throughout the proof, we use $C$ to denote a constant depending only on $d, p$, $q$ and $C_{1}$ and may change from line to line. We are going to apply the inequality (3.19) iteratively. For fixed $1 / 2 \leq \sigma^{\prime} \leq \sigma \leq 1$, and $k \in \mathbb{N}$ define

$$
\sigma_{k}=\sigma^{\prime}+2^{-k+1}\left(\sigma-\sigma^{\prime}\right)
$$

It is immediate that $\sigma_{k}-\sigma_{k+1}=2^{-k}\left(\sigma-\sigma^{\prime}\right)$ and that $\sigma_{1}=\sigma$, furthermore $\sigma_{k} \downarrow \sigma^{\prime}$. Let $p^{*}=p /(p-1)$. Then we have that $\rho>2 p^{*}$. Set $\alpha_{k}=\left(\rho / 2 p^{*}\right)^{k}, k \geq 1$. By definition, we have $\alpha_{k} \geq 1$ for $k \geq 1$. Let $\tilde{\eta}_{k}: \mathbb{R}^{d} \rightarrow \mathbb{R}$ be a smooth function which is identically 1 on $B\left(0, \sigma_{k+1} R\right)$ and vanishing on $\mathbb{R}^{d} \backslash B\left(0, \sigma_{k} R\right)$ and satisfies $\left\|\nabla \tilde{\eta}_{k}\right\| \leq \frac{2^{k}}{\left(\sigma-\sigma^{\prime}\right) R}$. Define a function $\eta_{k}$ by the restriction of $\tilde{\eta}_{k}$ to $W$. An application of (3.19) and of the relation $\alpha_{k} \rho=2 \alpha_{k+1} p^{*}$ yields

$$
\begin{aligned}
& \|u\|_{W\left(\sigma_{k+1} R\right), 2 \alpha_{k+1} p^{*}} \\
& \leq\left(C \frac{2^{2 k} \alpha_{k}^{2}\left|W\left(\sigma_{k} R\right)\right|^{2 / d}}{\left(\sigma-\sigma^{\prime}\right)^{2} R^{2}}\left\|\lambda^{-1}\right\|_{W\left(\sigma_{k} R\right), q}\|\Lambda\|_{W\left(\sigma_{k} R\right), p}\right)^{1 /\left(2 \alpha_{k}\right)}\|u\|_{W\left(\sigma_{k} R\right), 2 \alpha_{k} p^{*}}^{\gamma_{k}} \\
& \leq\left(C \frac{2^{2 k} \alpha_{k}^{2}}{\left(\sigma-\sigma^{\prime}\right)^{2}}\left\|\lambda^{-1}\right\|_{W(R), q}\|\Lambda\|_{W(R), p}\right)^{1 /\left(2 \alpha_{k}\right)}\|u\|_{W\left(\sigma_{k} R\right), 2 \alpha_{k} p^{*}}^{\gamma_{k}},
\end{aligned}
$$

where $\gamma_{k}=1$ if $\|u\|_{W\left(\sigma_{k} R\right), 2 \alpha_{k} p^{*}} \geq 1$ and $\gamma_{k}=1-1 / \alpha_{k}$ otherwise. Iterating the above inequality and stop at $k=1$, we get

$$
\|u\|_{W\left(\sigma_{j+1} R\right), 2 \alpha_{j+1} p^{*}} \leq \prod_{k=1}^{j}\left(C \frac{\left(\rho / p^{*}\right)^{2 k}}{\left(\sigma-\sigma^{\prime}\right)^{2}}\left\|\lambda^{-1}\right\|_{W(R), q}\|\Lambda\|_{W(R), p}\right)^{1 /\left(2 \alpha_{k}\right)}\|u\|_{W(\sigma R), \rho}^{\prod_{k=1}^{j} \gamma_{k}} .
$$

Observe that $\kappa:=\frac{1}{2} \sum \frac{1}{\alpha_{k}}<\infty, \sum \frac{k}{\alpha_{k}}<\infty$. Using (3.1), we have

$$
\begin{aligned}
& \|u\|_{W\left(\sigma^{\prime} R\right), 2 \alpha_{j} p^{*}} \leq\left(\frac{|W(\sigma R)|}{\left|W\left(\sigma^{\prime} R\right)\right|}\right)^{1 /\left(2 \alpha_{j} p^{*}\right)}\|u\|_{W(\sigma R), 2 \alpha_{j} p^{*}} \\
& \leq\left(\frac{C_{V}(\sigma R)^{d}}{\left(\sigma^{\prime} R\right)^{d}}\right)^{1 /\left(2 \alpha_{j} p^{*}\right)}\|u\|_{W(\sigma R), 2 \alpha_{j} p^{*}} \leq K\|u\|_{W(\sigma R), 2 \alpha_{j} p^{*}}
\end{aligned}
$$

for some $K>0$ and all $j \geq 1$. Hence, taking the limit as $j \rightarrow \infty$ gives the inequality

$$
\|u\|_{W\left(\sigma^{\prime} R\right), \infty} \leq C\left(\frac{1 \vee\left\|\lambda^{-1}\right\|_{W(R), q}\|\Lambda\|_{W(R), p}}{\left(\sigma-\sigma^{\prime}\right)^{2}}\right)^{\kappa}\|u\|_{W(\sigma R), \rho}^{\prod_{k=1}^{\infty} \gamma_{k}}
$$

Define $\gamma=\prod_{k=1}^{\infty}\left(1-1 / \alpha_{k}\right) \in(0,1]$, then $0<\gamma \leq \prod_{k=1}^{\infty} \gamma_{k}$ and the above inequality can be written as

$$
\|u\|_{W\left(\sigma^{\prime} R\right), \infty} \leq C\left(\frac{1 \vee\left\|\lambda^{-1}\right\|_{W(R), q}\|\Lambda\|_{W(R), p}}{\left(\sigma-\sigma^{\prime}\right)^{2}}\right)^{\kappa}\|u\|_{W(\sigma R), \rho}^{\gamma} \vee\|u\|_{W(\sigma R), \rho}
$$

which is the desired inequality.
The main goal of this section is the following inequality. It is proved as in [9, Corollary 2.2].

Corollary 3.9 (maximal inequality). Let $R \geq R_{I} \vee R_{V}$. Suppose that $u$ is a solution of (3.17) in $W(R)$. Then, for all $\alpha \in(0, \infty)$ and for any $1 / 2 \leq \sigma^{\prime}<\sigma<1$ there exist $C^{\prime}=C^{\prime}\left(p, q, d, c_{f}\right)>0, \gamma^{\prime}=\gamma^{\prime}(\gamma, \alpha, \rho)$ and $\kappa^{\prime}=\kappa^{\prime}(\kappa, \alpha, \rho)$, such that
(3.21) $\|u\|_{W\left(\sigma^{\prime} R\right), \infty} \leq C^{\prime}\left(\frac{1 \vee\left\|\lambda^{-1}\right\|_{W(R), q}\|\Lambda\|_{W(R), p}}{\left(\sigma-\sigma^{\prime}\right)^{2}}\right)^{\kappa^{\prime}}\|u\|_{W(\sigma R), \alpha}^{\gamma^{\prime}} \vee\|u\|_{W(\sigma R), \alpha}$.

## 4. Sublinearity

In this section, we will prove the sublinearity of the corrector and non-degeneracy of the covariance matrix. First, we will prove the $L^{p}$-sublinearity. In [9], the proof relies on the fact that $\mathbb{E}\left[U^{k}\right]=0$. However, as we mentioned in Remark 2.2, since we consider other measure $\widehat{\mathbb{P}}$, we have $\widehat{\mathbb{E}}\left[U^{k}\right] \neq 0$. To overcome this issue, we will take another approach.

Lemma 4.1. Let $\varphi \in \mathcal{C}$ and define $\Phi(x, \omega)$ by

$$
\Phi(x, \omega)=\int_{0}^{1}\left|x \cdot \nabla \varphi\left(\tau_{t x} \omega\right)\right| d t
$$

Then,

$$
\sup _{x \in \mathbb{R}^{d}} \Phi(x, \omega)<\infty
$$

holds $\widehat{\mathbb{P}}$-almost surely $\omega$.

Proof. Let

$$
\varphi(\omega)=\int_{\mathbb{R}^{d}} f\left(\tau_{x} \omega\right) \eta(x) d x, \quad \eta \in C_{c}^{\infty}\left(\mathbb{R}^{d}\right)
$$

and set $K=\operatorname{supp} \eta$. Denote the diameter of $K$ by $\delta$. Then we have

$$
\begin{aligned}
\Phi(x, \omega) & =\int_{0}^{1}\left|\left\langle x, \nabla \varphi\left(\tau_{t x} \omega\right)\right\rangle\right| d t \\
& \leq \int_{0}^{1} \int_{\mathbb{R}^{d}}\left|\left\langle f\left(\tau_{t x+y} \omega\right) x, \nabla \eta(y)\right\rangle\right| d y d t \\
& =\int_{0}^{1} \int_{K}\left|\left\langle f\left(\tau_{t x+y} \omega\right) x, \nabla \eta(y)\right\rangle\right| d y d t \\
& \leq \int_{0}^{1} \int_{K} \delta\|f\|_{\infty}\|\nabla \eta\|_{\infty} d y d t \\
& \leq \delta|K|\|f\|_{\infty}\|\nabla \eta\|_{\infty}
\end{aligned}
$$

Hence we get the desired result.
Set $\chi_{\varepsilon}(x, \omega)=\varepsilon \chi(x / \varepsilon, \omega)$. To prove the $L^{p}$-sublinearity, we recall a functional analysis result. Let $\mathscr{B}_{1}, \mathscr{B}_{2}$ be Banach spaces and $T: \mathscr{B}_{1} \rightarrow \mathscr{B}_{2}$ a compact operator. Then, for each sequence $\left\{x_{n}\right\}_{n} \subset \mathscr{B}_{1}$ such that $x_{n} \rightarrow x$ weakly, we have that $T x_{n} \rightarrow T x$ strongly.

Lemma 4.2 ( $L^{p}$-SUblinearity). For $\widehat{\mathbb{P}}$-a.s. $\omega$ and $R \geq R_{0}(\omega)$,

$$
\lim _{\varepsilon \rightarrow 0}\left\|\chi_{\varepsilon}^{k}(x, \omega)\right\|_{2 p^{*}, W_{R}}=0
$$

for $k=1, \ldots, d$.
Proof. It is enough to show that for any $\eta \in C_{c}^{\infty}\left(W_{R}^{\prime}\right)$ we have

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} \frac{1}{\left|W_{R}^{\prime}\right|} \int_{W_{R}^{\prime}} \chi_{\varepsilon}(x, \omega) \eta(x) d x=0 \tag{4.1}
\end{equation*}
$$

Indeed, the above property implies the weak convergence $\chi_{\varepsilon} \rightarrow 0$ in $L^{2}\left(W_{R}^{\prime}\right)$. This gives the strong convergence in $L^{2 p^{*}}\left(W_{R}^{\prime}\right)$, because $W^{1, \frac{2 q}{q+1}}\left(W_{R}^{\prime}\right)$ is compactly embedded in $L^{2 p^{*}}\left(W_{R}^{\prime}\right)$ and the sequence $\left\{\chi_{\varepsilon}\right\}_{\varepsilon>0}$ is bounded in $W^{1, \frac{2 q}{q+1}}\left(W_{R}^{\prime}\right)$. To show the equality (4.1), we will be apart from [9]. The following argument is motivated by [10].

We have

$$
\begin{aligned}
& \left|\frac{1}{\left|W_{R}^{\prime}\right|} \int_{W_{R}^{\prime}} \chi_{\varepsilon}(x, \omega) \eta(x) d x\right| \\
& =\left|\frac{\varepsilon^{d+1}}{\left|W_{R}^{\prime}\right|} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}} \chi(z, \omega) \eta(\varepsilon z) d z\right| \\
& \leq \frac{\varepsilon}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}}|\chi(z, \omega) \eta(\varepsilon z)| d z
\end{aligned}
$$

$$
\begin{aligned}
& \leq \frac{\varepsilon}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}} \int_{0}^{1}\left|z \cdot U^{k}\left(\tau_{t z} \omega\right) \eta(\varepsilon z)\right| d t d z \\
& \leq\|\eta\|_{\infty} \frac{\varepsilon}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{0}^{1} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}}\left|z \cdot U^{k}\left(\tau_{t z} \omega\right)\right| d z d t
\end{aligned}
$$

Now, let $\varphi_{n} \in \mathcal{C}$ be an approximate sequence such that $\nabla \varphi_{n} \rightarrow U^{k}$ in $\widehat{\mathbb{L}}^{2}$ and set

$$
\Phi_{n}(x, \omega)=\int_{0}^{1}\left|x \cdot \varphi_{n}\left(\tau_{t x} \omega\right)\right| d t
$$

Then we compute

$$
\begin{aligned}
& \frac{\varepsilon}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{0}^{1} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}}\left|z \cdot U^{k}\left(\tau_{t z} \omega\right)\right| d z d t \\
& \leq \frac{\varepsilon}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{0}^{1} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}}\left|z \cdot \nabla \varphi_{n}\left(\tau_{t z} \omega\right)\right| d z d t \\
& +\frac{\varepsilon}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{0}^{1} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}}\left|z \cdot U^{k}\left(\tau_{t z} \omega\right)-z \cdot \nabla \varphi_{n}\left(\tau_{t z} \omega\right)\right| d z d t
\end{aligned}
$$

The first term of the right hand side is bounded above by

$$
\begin{aligned}
& \frac{\varepsilon}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{0}^{1} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}}\left|z \cdot \nabla \varphi_{n}\left(\tau_{t z} \omega\right)\right| d z d t \\
& =\frac{\varepsilon}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}} \int_{0}^{1}\left|z \cdot \nabla \varphi_{n}\left(\tau_{t z} \omega\right)\right| d t d z \\
& =\frac{\varepsilon}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}} \Phi_{n}(z, \omega) d t d z \\
& \leq \varepsilon \sup _{z \in \mathbb{R}^{d}} \Phi_{n}(z, \omega) .
\end{aligned}
$$

Since $\sup _{z \in \mathbb{R}^{d}}\left|\Phi_{n}(z, \omega)\right|<\infty$ by Lemma 4.1, this term tends to zero as $\varepsilon \rightarrow 0$. For the second term, using the Cauchy-Schwarz inequality, we estimate

$$
\begin{aligned}
& \frac{\varepsilon}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{0}^{1} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}}\left|z \cdot U^{k}\left(\tau_{t z} \omega\right)-z \cdot \nabla \varphi_{n}\left(\tau_{t z} \omega\right)\right| d z d t \\
& =\frac{1}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{0}^{1} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}}\left|(\varepsilon z) \cdot\left(U^{k}\left(\tau_{t z} \omega\right)-\nabla \varphi_{n}\left(\tau_{t z} \omega\right)\right)\right| d z d t \\
& \leq \frac{1}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{0}^{1} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}} \sqrt{|\varepsilon z|\left|U^{k}\left(\tau_{t z} \omega\right)-\nabla \varphi_{n}\left(\tau_{t z} \omega\right)\right|} d z d t
\end{aligned}
$$

Using the fact that $|z| \leq R / \varepsilon$ for $z \in \frac{1}{\varepsilon} W_{R}$ and the change of variables, we have

$$
\begin{aligned}
& \frac{1}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{0}^{1} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}} \sqrt{|\varepsilon z|\left|U^{k}\left(\tau_{t z} \omega\right)-\nabla \varphi_{n}\left(\tau_{t z} \omega\right)\right|} d z d t \\
& \leq R^{\frac{1}{2}} \frac{1}{\left|\frac{1}{\varepsilon} W_{R}^{\prime}\right|} \int_{0}^{1} \int_{\frac{1}{\varepsilon} W_{R}^{\prime}} \sqrt{\left|U^{k}\left(\tau_{t z} \omega\right)-\nabla \varphi_{n}\left(\tau_{t z} \omega\right)\right|} d z d t
\end{aligned}
$$

$$
=R^{\frac{1}{2}} \int_{0}^{1} \frac{1}{\left|\frac{t}{\varepsilon} W_{R}^{\prime}\right|} \int_{\frac{t}{\varepsilon} W_{R}^{\prime}} \sqrt{\left|U^{k}\left(\tau_{u} \omega\right)-\nabla \varphi_{n}\left(\tau_{u} \omega\right)\right|} d u d t
$$

By the volume regularity ((1) of Assumption 4), this is bounded above by

$$
C R^{\frac{1}{2}} \int_{0}^{1} \frac{1}{\left|B\left(0, \frac{t R}{\varepsilon}\right)\right|} \int_{B\left(0, \frac{t R}{\varepsilon}\right)} \sqrt{\left|U^{k}\left(\tau_{u} \omega\right)-\nabla \varphi_{n}\left(\tau_{u} \omega\right)\right|} 1_{\left\{0 \in W^{\prime}\left(\tau_{u} \omega\right)\right\}} d u d t
$$

By the ergodic theorem, we have

$$
\lim _{\varepsilon \rightarrow 0} \frac{1}{\left|B\left(0, \frac{t R}{\varepsilon}\right)\right|} \int_{B\left(0, \frac{t R}{\varepsilon}\right)} \sqrt{\left|U^{k}\left(\tau_{u} \omega\right)-\nabla \varphi_{n}\left(\tau_{u} \omega\right)\right|} 1_{\left\{0 \in W^{\prime}\left(\tau_{u} \omega\right)\right\}} d u=\widehat{\mathbb{E}}\left[\sqrt{\left|U^{k}-\nabla \varphi_{n}\right|}\right]
$$

Letting $n$ tend to infinity, this term converges to zero. Now we have

$$
\begin{aligned}
& \left|\frac{1}{\left|W_{R}^{\prime}\right|} \int_{\mathbb{R}^{d}} \chi_{\varepsilon}(x, \omega) \eta(x) d x\right| \\
& \leq\|\eta\|_{\infty}\left(\varepsilon \sup _{z \in \mathbb{R}^{d}}\left|\Phi_{n}(z, \omega)\right|\right. \\
& \left.+C_{V}^{-1} R \int_{0}^{1} \frac{1}{\left|B\left(0, \frac{t R}{\varepsilon}\right)\right|} \int_{B\left(0, \frac{t R}{\varepsilon}\right)}\left|U^{k}\left(\tau_{u} \omega\right)-\nabla \varphi_{n}\left(\tau_{u} \omega\right)\right| 1_{\left\{0 \in W^{\prime}\left(\tau_{u} \omega\right)\right\}} d u d t\right) .
\end{aligned}
$$

First let $\varepsilon$ tend to zero and then $n$ tend to infinity, we get (4.1) and obtain the result.
Proposition 4.3 ( $L^{\infty}$-sublinearity). For all $R>0$,

$$
\lim _{\varepsilon \rightarrow 0} \sup _{x \in W_{R}^{\prime}}\left|\chi_{\varepsilon}(x, \omega)\right|=0, \quad \widehat{\mathbb{P}} \text {-a.s. }
$$

Proof. Observe that by Proposition 2.8, the function $\chi_{\varepsilon}^{k}(x, \omega)$ is a solution of

$$
\mathcal{E}^{\omega}\left(\chi_{\varepsilon}, \varphi\right)=\int_{W^{\prime}}\left\langle a\left(\tau_{x} \omega\right) \nabla f_{k}, \nabla \varphi\right\rangle d x
$$

in $W_{R}^{\prime}$ for all $\varepsilon>0$. Here $f_{k}(x)=x_{k}$ and $\varphi \in H^{1}\left(W_{R}^{\prime}\right) \cap C_{c}^{\infty}\left(\overline{W_{R}}\right)$. We first consider the case $R \geq R_{0}(\omega)$. Since $\left|\nabla f_{k}\right| \leq 1$, by Lemma 4.2 we have

$$
\lim _{\varepsilon \rightarrow 0}\left\|\chi_{\varepsilon}^{k}(x, \omega)\right\|_{2 p^{*}, W_{R}^{\prime}}=0
$$

Therefore, for $R \geq R_{0}(\omega)$, we get the desired result by the maximal inequality (3.21) with $\alpha=2 p^{*}$. It remains to treat the case that $R \in\left(0, R_{0}(\omega)\right)$, but in that case, the desired bound immediately follows from the fact that $\sup _{x \in W_{R}^{\prime}}\left|\chi_{\varepsilon}(x, \omega)\right| \leq \sup _{x \in W_{R_{0}}^{\prime}}\left|\chi_{\varepsilon}(x, \omega)\right|$ for $R \in\left(0, R_{0}\right)$.

To prove the positive-definiteness of the covariance matrix, we show the following lemma. Recall that $T_{x}$ is the operator on $L^{\infty}(\widehat{\mathbb{P}})$ defined by $T_{x} G(\omega)=G\left(\tau_{x} \omega\right)$.

Lemma 4.4. Let $G: \Omega \rightarrow \mathbb{R}$ be an integrable random variable such that $G(\omega)=0$
for $\widehat{\mathbb{P}}$-almost all $\omega \in \widehat{\Omega}$. Fix $x \in \mathbb{R}^{d}$. Then for $\widehat{\mathbb{P}}$-almost all $\omega$,

$$
\int_{0}^{1}\left(T_{\gamma(t)} G\right) 1_{\left\{\gamma(t) \in W^{\prime}(\omega)\right\}} \gamma^{\prime}(t) d t=0
$$

holds for all smooth path $\gamma:[0,1] \rightarrow \mathbb{R}^{d}$ satisfying $\gamma(0)=0$ and $\gamma(1)=x$.
Proof. First we fix a smooth path $\gamma$. Then for all $F \in L^{\infty}(\widehat{\mathbb{P}})$, by Fubini's theorem and the definition of $\widehat{\mathbb{P}}$, we have

$$
\begin{aligned}
& \widehat{\mathbb{E}}\left[F \int_{0}^{1}\left(T_{\gamma(t)} G\right) 1_{\left\{\gamma(t) \in W^{\prime}\right\}} \gamma^{\prime}(t) d t\right] \\
& =\int_{0}^{1} \widehat{\mathbb{E}}\left[F\left(T_{\gamma(t)} G\right) 1_{\left\{\gamma(t) \in W^{\prime}\right\}} \gamma^{\prime}(t)\right] d t \\
& =\int_{0}^{1} \mathbb{E}\left[F\left(T_{\gamma(t)} G\right) 1_{\left\{\gamma(t) \in W^{\prime}\right\}} \gamma^{\prime}(t) 1_{\left\{0 \in W^{\prime}\right\}}\right] \mathbb{P}\left(0 \in W^{\prime}\right)^{-1} d t
\end{aligned}
$$

Using the translation invariance of $\mathbb{P}$, the last line equals to

$$
\begin{aligned}
& \int_{0}^{1} \mathbb{E}\left[\left(T_{-\gamma(t)} F\right) G 1_{\left\{0 \in W^{\prime}\right\}} 1_{\left\{-\gamma(t) \in W^{\prime}\right\}} \gamma^{\prime}(t)\right] \mathbb{P}\left(0 \in W^{\prime}\right)^{-1} d t \\
& =\int_{0}^{1} \widehat{\mathbb{E}}\left[\left(T_{-\gamma(t)} F\right) G \gamma^{\prime}(t) 1_{\left\{-\gamma(t) \in W^{\prime}\right\}}\right] d t \\
& =\int_{0}^{1} \widehat{\mathbb{E}}\left[\left(T_{-\gamma(t)} F\right) \cdot 0 \cdot \gamma^{\prime}(t) 1_{\left\{-\gamma(t) \in W^{\prime}\right\}}\right] d t \\
& =0
\end{aligned}
$$

Since $F \in L^{\infty}(\widehat{\mathbb{P}})$ is arbitrary, we obtain

$$
\int_{0}^{1}\left\langle T_{\gamma(t)} G, \gamma^{\prime}(t)\right\rangle 1_{\left\{\gamma(t) \in W^{\prime}\right\}} d t=0
$$

for $\widehat{\mathbb{P}}$-a.s. $\omega$. This implies the result because the collection of all smooth path is separable with respect to the supremum norm (that is, for a path $\gamma:[0,1] \rightarrow \mathbb{R}^{d}$, define its norm by $\left.\sup _{t \in[0,1]}|\gamma(t)|\right)$.

Proposition 4.5. The covariance matrix $\mathbb{D}$ is positive-definite.
Proof. The general idea of the proof goes back to [10]. Assume that there is a vector $v \in \mathbb{R}^{d}$ with $|v|=1$ such that $\langle v, \mathbb{D} v\rangle=0$. We have

$$
\begin{aligned}
\langle v, \mathbb{D} v\rangle & =\sum_{i=1}^{d} v_{i}(\mathbb{D} v)_{i} \\
& =\sum_{i=1}^{d} v_{i}\left(\sum_{j=1}^{d} d_{i j} v_{j}\right) \\
& =\sum_{k=1} \widehat{\mathbb{E}}[\langle a \xi, \xi\rangle]
\end{aligned}
$$

where $\xi=\left(\left\langle v, D_{1} y\right\rangle, \ldots,\left\langle v, D_{d} y\right\rangle\right), y=\left(y^{1}, \ldots, y^{d}\right)$ and $D_{k} y=\left(D_{k} y^{1}, \ldots, D_{k} y^{d}\right)$. Thus, we have $\left\langle v, D_{k} y(0, \omega)\right\rangle=0$ for $k=1, \ldots, d \widehat{\mathbb{P}}$-a.s. Take $\omega \in \widehat{\Omega}$ and smooth path $\gamma_{\omega}:[0,1] \rightarrow \mathbb{R}^{d}$ such that $\gamma_{\omega}(0)=0, \gamma_{\omega}(1)=x$ and $\gamma_{\omega}([0,1]) \subset W^{\prime}(\omega)$. We denote the $k$-th component of $\gamma_{\omega}^{\prime}$ by $\gamma_{\omega, k}^{\prime}$. Since $D_{k} y(x, \omega)=D_{k} y\left(0, \tau_{x} \omega\right)$ and

$$
y^{i}(x, \omega)=\int_{0}^{1}\left\langle\nabla y^{i}\left(\gamma_{\omega}(t), \omega\right), \gamma_{\omega}^{\prime}(t)\right\rangle d t=\sum_{k=1}^{d} \int_{0}^{1} D_{k} y^{i}\left(0, \tau_{\gamma_{\omega}(t)} \omega\right) \gamma_{\omega, k}^{\prime}(t) 1_{\left\{\gamma_{\omega}(t) \in W^{\prime}(\omega)\right\}} d t
$$

it holds that

$$
\begin{aligned}
\langle v, y(x, \omega)\rangle & =\sum_{i=1}^{d} \sum_{k=1}^{d} \int_{0}^{1} v_{i} D_{k} y^{i}\left(0, \tau_{\gamma(t)} \omega\right) \gamma_{\omega, k}^{\prime}(t) 1_{\left\{\gamma_{\omega}(t) \in W^{\prime}(\omega)\right\}} d t \\
& =\sum_{k=1}^{d} \int_{0}^{1}\left\langle v, D_{k} y\left(0, \tau_{\gamma_{\omega}(t)} \omega\right)\right\rangle 1_{\left\{\gamma_{\omega}(t) \in W^{\prime}\right\}} \gamma_{\omega}^{\prime}(t) d t .
\end{aligned}
$$

Thus it follows from Lemma 4.4 that we have $\langle v, y(x, \omega)\rangle=0$. Hence $\langle v, x\rangle=\langle v, \chi(x, \omega)\rangle$ $\widehat{\mathbb{P}}$-a.s. However, it implies

$$
0=\lim _{\varepsilon \rightarrow 0} \sup _{x \in W_{R}^{\prime}}\langle v, \varepsilon \chi(x / \varepsilon, \omega)\rangle=\lim _{\varepsilon \rightarrow 0} \sup _{x \in W_{R}^{\prime}}\langle v, \varepsilon(x / \varepsilon)\rangle=\lim _{\varepsilon \rightarrow 0} \sup _{x \in W_{R}^{\prime}}\langle v, x\rangle>0,
$$

a contradiction.

## 5. Proof of the main theorem

Recall that we have the decomposition

$$
\varepsilon X_{\cdot / \varepsilon^{2}}^{\omega}=\varepsilon \chi\left(X_{\cdot / \varepsilon^{2}}^{\omega}, \omega\right)+\varepsilon y\left(X_{\cdot / \varepsilon^{2}}^{\omega}, \omega\right) .
$$

First, we show that the martingale $\varepsilon y\left(X_{\cdot / \varepsilon^{2}}^{\omega}, \omega\right)$ converges to a Brownian motion with covariance matrix $\mathbb{D}$.

Lemma 5.1. Set $M_{t}^{\omega}=y\left(X_{t}^{\omega}, \omega\right)$. Then $\varepsilon M_{\cdot / \varepsilon^{2}}^{\omega}$ converges to a Brownian motion with the covariance matrix $\mathbb{D}$.

Proof. By Corollary 2.10 and the ergodic theorem for the environment process $\left\{\tau_{X_{t} \omega} \omega\right\}_{t}$ (see [26, Proposition 2.1]), we have

$$
\begin{aligned}
\left\langle M_{\cdot}^{k, \varepsilon}, M_{\cdot}^{\ell, \varepsilon}\right\rangle_{t} & =\varepsilon^{2} \int_{0}^{t / \varepsilon^{2}} \sum_{i=1}^{d}\left\langle a\left(\tau_{X_{s}^{\omega}} \omega\right) \nabla y^{k}\left(X_{s}, \omega\right), \nabla y^{\ell}\left(X_{s}, \omega\right)\right\rangle d s \\
& \rightarrow \widehat{\mathbb{E}}\left[\left\langle a(\omega) \nabla y^{k}(0, \omega), \nabla y^{\ell}(0, \omega)\right\rangle\right] .
\end{aligned}
$$

Hence, we get the result by [20, Theorem 5.1.].
It remains to show that the corrector $\varepsilon \chi\left(X_{t / \varepsilon^{2}}^{\omega}, \omega\right)$ converges to zero in distribution. For that the sublinearity of the corrector will play a major role.

Let $T>0$ be a fixed time horizon. We claim that for all $\delta>0$

$$
\begin{equation*}
\lim _{\varepsilon \rightarrow 0} P_{0}^{\omega}\left(\sup _{0 \leq t \leq T}\left|\varepsilon \chi\left(X_{t / \varepsilon^{2}}^{\omega}, \omega\right)\right|>\delta\right)=0 \tag{5.1}
\end{equation*}
$$

Denote by $\tau_{R}^{\varepsilon, \omega}$ the exit time of $\varepsilon X_{t / \varepsilon^{2}}^{\omega}$ from $W_{R}^{\prime}$. Observe that
(5.2) $\quad \limsup _{\varepsilon \rightarrow 0} P_{0}^{\omega}\left(\sup _{0 \leq t \leq T}\left|\varepsilon \chi\left(X_{t / \varepsilon^{2}}^{\omega}, \omega\right)\right|>\delta\right)$
$\leq \limsup _{\varepsilon \rightarrow 0} P_{0}^{\omega}\left(\left|\sup _{0 \leq t \leq \tau_{R}^{\varepsilon, \omega}} \varepsilon \chi\left(X_{t / \varepsilon^{2}}^{\omega}, \omega\right)\right|>\delta\right)+\limsup _{\varepsilon \rightarrow 0} P_{0}^{\omega}\left(\left|\sup _{0 \leq t \leq T} \varepsilon X_{t / \varepsilon^{2}}^{\omega}\right|>R\right)$.
By Proposition 4.3, we have

$$
\lim _{\varepsilon \rightarrow 0} \sup _{0 \leq t \leq \tau_{R}^{\varepsilon, \omega}}\left|\varepsilon \chi\left(X_{t / \varepsilon^{2}}^{\omega}, \omega\right)\right|=0
$$

and therefore $\widehat{\mathbb{P}}$-almost surely

$$
\limsup _{\varepsilon \rightarrow 0} P_{0}^{\omega}\left(\sup _{0 \leq t \leq \tau_{R}^{\varepsilon, \omega}}\left|\varepsilon \chi\left(X_{t / \varepsilon^{2}}^{\omega}, \omega\right)\right|>\delta\right)=0
$$

To obtain a bound of the second term of (5.2), we use again Proposition 4.3 to say that there exists $\tilde{\varepsilon}(\omega)>0$, which may depend on $\omega$ such that for all $\varepsilon<\tilde{\varepsilon}$ we have $\sup _{0 \leq t \leq \tau_{R}^{\varepsilon, \omega}}\left|\varepsilon \chi\left(X_{t / \varepsilon^{2}}^{\omega}, \omega\right)\right|<1$. For such $\varepsilon$ we have $\widehat{\mathbb{P}}$-almost surely

$$
\begin{aligned}
P_{0}^{\omega}\left(\sup _{0 \leq t \leq T}\left|\varepsilon X_{t / \varepsilon^{2}}^{\omega}\right| \geq R\right) & =P_{0}^{\omega}\left(\tau_{R}^{\varepsilon, \omega} \leq T\right) \\
& =P_{0}^{\omega}\left(\tau_{R}^{\varepsilon, \omega} \leq T, \sup _{0 \leq t \leq \tau_{R}^{\varepsilon, \omega}}\left|\varepsilon y\left(X_{t / \varepsilon^{2}}^{\omega}, \omega\right)\right|>R-1\right) \\
& \leq P_{0}^{\omega}\left(\sup _{0 \leq t \leq T}\left|\varepsilon y\left(X_{t / \varepsilon^{2}}^{\omega}, \omega\right)\right|>R-1\right)
\end{aligned}
$$

Thanks to Lemma 5.1, the process $\varepsilon y\left(X_{\cdot / \varepsilon^{2}}^{\omega}, \omega\right)$ converges in distribution under $P_{0}^{\omega}$ to a non-degenerate Brownian motion with the deterministic covariance matrix given by $\mathbb{D}$. Hence there exist positive constants $c_{1}, c_{2}$ independent of $\varepsilon$ and $\omega$ such that

$$
\limsup _{\varepsilon \rightarrow 0} P_{0}^{\omega}\left(\sup _{0 \leq t \leq T}\left|\varepsilon y\left(X_{t / \varepsilon^{2}}^{\omega}, \omega\right)\right|>R-1\right) \leq c_{1} e^{-c_{2} R}
$$

from which it follows

$$
\limsup _{\varepsilon \rightarrow 0} P_{0}^{\omega}\left(\sup _{0 \leq t \leq T}\left|\varepsilon X_{t / \varepsilon^{2}}^{\omega}\right|>R\right) \leq c_{1} e^{-c_{2} R}
$$

Therefore,

$$
\limsup _{\varepsilon \rightarrow 0} P_{0}^{\omega}\left(\sup _{0 \leq t \leq T}\left|\varepsilon \chi\left(X_{t / \varepsilon^{2}}^{\omega}, \omega\right)\right|>\delta\right) \leq c_{1} e^{-c_{2} R}
$$

and since $R>1$ was arbitrary, the claim (5.1) follows, namely the corrector converges to zero in law under $P_{0}^{\omega}, \widehat{\mathbb{P}}$-almost surely.

The convergence to zero in law of the corrector $\varepsilon \chi\left(X_{\cdot / \varepsilon^{2}}^{\omega}, \omega\right)$, combined with the fact that $\varepsilon y\left(X_{\cdot / \varepsilon^{2}}^{\omega}, \omega\right)$ satisfies an invariance principle $\widehat{\mathbb{P}}$-almost surely and that $\varepsilon X_{\cdot / \varepsilon^{2}}^{\omega}=$ $\varepsilon \chi\left(X_{\cdot / \varepsilon^{2}}^{\omega}, \omega\right)+\varepsilon y\left(X_{\cdot / \varepsilon^{2}}^{\omega}, \omega\right)$, implies that the family $\varepsilon X_{\cdot / \varepsilon^{2}}^{\omega}$ under $P_{0}^{\omega}$ satisfies an invariance principle $\widehat{\mathbb{P}}$-almost surely with the same limiting law.

Acknowledgements The author would like to thank Professor Hideki Tanemura for suggesting the problem and fruitful discussions. This work was supported by Keio University Doctoral Student Get-in-Aid Program.

## References

[1] S. Andres. Invariance principle for the random conductance model with dynamic bounded conductances. Ann. Inst. Henri Poincaré Probab. Stat., 50(2):352-374, 2014.
[2] S. Andres, M. T. Barlow, J.-D. Deuschel, and B. M. Hambly. Invariance principle for the random conductance model. Probab. Theory Related Fields, 156(3-4):535-580, 2013.
[3] S. Andres, J.-D. Deuschel, and M. Slowik. Invariance principle for the random conductance model in a degenerate ergodic environment. Ann. Probab., 43(4):1866-1891, 2015.
[4] M. T. Barlow. Random walks on supercritical percolation clusters. Ann. Probab., 32(4):3024-3084, 2004.
[5] M. T. Barlow and J.-D. Deuschel. Invariance principle for the random conductance model with unbounded conductances. Ann. Probab., 38(1):234-276, 2010.
[6] R. F. Bass and P. Hsu. Some potential theory for reflecting Brownian motion in Hölder and Lipschitz domains. Ann. Probab., 19(2):486-508, 1991.
[7] N. Berger and M. Biskup. Quenched invariance principle for simple random walk on percolation clusters. Probab. Theory Related Fields, 137(1-2):83-120, 2007.
[8] M. Biskup and T. M. Prescott. Functional CLT for random walk among bounded random conductances. Electron. J. Probab., 12:no. 49, 1323-1348, 2007.
[9] A. Chiarini and J.-D. Deuschel. Invariance principle for symmetric diffusions in a degenerate and unbounded stationary and ergodic random medium. Ann. Inst. Henri Poincaré Probab. Stat., 52(4):1535-1563, 2016.
[10] J.-D. Deuschel, T. A. Nguyen, and M. Slowik. Quenched invariance principles for the random conductance model on a random graph with degenerate ergodic weights. Probab. Theory Related Fields, 170(1-2):363-386, 2018.
[11] L. C. Evans and R. F. Gariepy. Measure theory and fine properties of functions. Textbooks in Mathematics. CRC Press, Boca Raton, FL, revised edition, 2015.
[12] A. Fannjiang and T. Komorowski. Correction: "an invariance principle for diffusion in turbulence" [ann. probab. 27 (1999), no. 2, 751-781; mr1698963 (2001e:60069)]. Ann. Probab., 30(1):480-482.
[13] A. Fannjiang and T. Komorowski. A martingale approach to homogenization of unbounded random flows. Ann. Probab., 25(4):1872-1894, 1997.
[14] A. Fannjiang and T. Komorowski. An invariance principle for diffusion in turbulence. Ann. Probab., 27(2):751-781, 1999.
[15] M. Fukushima, S. Nakao, and M. Takeda. On Dirichlet forms with random data-recurrence and homogenization. In Stochastic processes-mathematics and physics, II (Bielefeld, 1985), volume 1250 of Lecture Notes in Math., pages 87-97. Springer, Berlin, 1987.
[16] M. Fukushima, Y. Oshima, and M. Takeda. Dirichlet forms and symmetric Markov processes, volume 19 of De Gruyter Studies in Mathematics. Walter de Gruyter \& Co., Berlin, extended edition, 2011.
[17] M. Fukushima and M. Tomisaki. Reflecting diffusions on Lipschitz domains with cusps-analytic construction and Skorohod representation. volume 4, pages 377-408. 1995. Potential theory and degenerate partial differential operators (Parma).
[18] M. Fukushima and M. Tomisaki. Construction and decomposition of reflecting diffusions on Lipschitz domains with Hölder cusps. Probab. Theory Related Fields, 106(4):521-557, 1996.
[19] D. Gilbarg and N. S. Trudinger. Elliptic partial differential equations of second order. Classics in Mathematics. Springer-Verlag, Berlin, 2001. Reprint of the 1998 edition.
[20] I. S. Helland. Central limit theorems for martingales with discrete or continuous time. Scand. J. Statist., 9(2):79-94, 1982.
[21] C. Kipnis and S. R. S. Varadhan. Central limit theorem for additive functionals of reversible Markov processes and applications to simple exclusions. Comm. Math. Phys., 104(1):1-19, 1986.
[22] P. Mathieu. Quenched invariance principles for random walks with random conductances. J. Stat. Phys., 130(5):1025-1046, 2008.
[23] P. Mathieu and A. Piatnitski. Quenched invariance principles for random walks on percolation clusters. Proc. R. Soc. Lond. Ser. A Math. Phys. Eng. Sci., 463(2085):2287-2307, 2007.
[24] K. Matsuura. Doubly Feller property of Brownian motions with Robin boundary condition. Potential Anal., 53(1):23-53, 2020.
[25] R. Meester and R. Roy. Continuum percolation, volume 119 of Cambridge Tracts in Mathematics. Cambridge University Press, Cambridge, 1996.
[26] H. Osada. Homogenization of diffusion processes with random stationary coefficients. In Probability theory and mathematical statistics (Tbilisi, 1982), volume 1021 of Lecture Notes in Math., pages 507-517. Springer, Berlin, 1983.
[27] H. Osada. Homogenization of reflecting barrier Brownian motions. In Asymptotic problems in probability theory: stochastic models and diffusions on fractals (Sanda/Kyoto, 1990), volume 283 of Pitman Res. Notes Math. Ser., pages 59-74. Longman Sci. Tech., Harlow, 1993.
[28] H. Osada and T. Saitoh. An invariance principle for non-symmetric Markov processes and reflecting diffusions in random domains. Probab. Theory Related Fields, 101(1):45-63, 1995.
[29] R. Rhodes. On homogenization of space-time dependent and degenerate random flows. Stochastic Process. Appl., 117(10):1561-1585, 2007.
[30] R. Rhodes. On homogenization of space-time dependent and degenerate random flows. II. Ann. Inst. Henri Poincaré Probab. Stat., 44(4):673-692, 2008.
[31] V. Sidoravicius and A.-S. Sznitman. Quenched invariance principles for walks on clusters of percolation or among random conductances. Probab. Theory Related Fields, 129(2):219-244, 2004.
[32] A.-S. Sznitman and O. Zeitouni. An invariance principle for isotropic diffusions in random environment. Invent. Math., 164(3):455-567, 2006.
[33] H. Tanemura. Behavior of the supercritical phase of a continuum percolation model on $\mathbf{R}^{d}$. J. Appl. Probab., 30(2):382-396, 1993.
[34] H. Tanemura. Homogenization of a reflecting barrier Brownian motion in a continuum percolation cluster in $\mathbf{R}^{d}$. Kodai Math. J., 17(2):228-245, 1994.

Yutaka TAKEUCHI
Department of Mathematics, Faculty of Science and Technology, Keio University 3-14-1 Hiyoshi, Kohoku-ku, Yokohama, 223-8522, Japan
E-mail: yutaka.takeuchi@keio.jp


[^0]:    1991 Mathematics Subject Classification.
    Key Words and Phrases. Continuum percolation, Quenched invariance principle.

